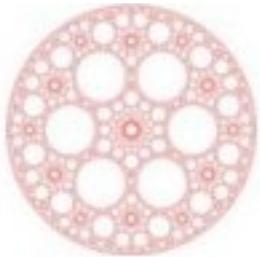


# Recovering Entanglement by Local Operations

Giuliano Benenti



Center for Nonlinear and Complex Systems  
Univ. Insubria, Como, Italy

In collaboration with:

Antonio D'Arrigo, Elisabetta Paladino, Giuseppe Falci (Catania)  
Rosario Lo Franco (Palermo)

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# Outline

Problem: entanglement revivals without non-local operations

Quantification of entanglement based on the density operator formalism vs ensemble description

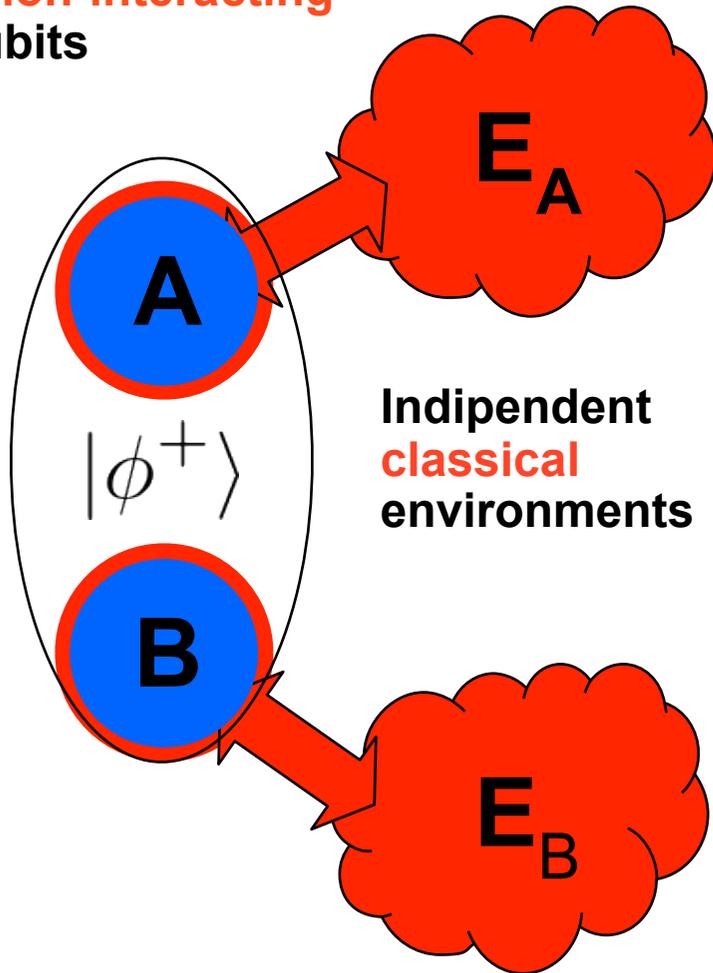
Definition of hidden entanglement

Examples: Random local fields, solid state system affected by low frequency noise, amplitude damping channel, etc.

Non-Markovian dynamics and non-monotonous behavior of entanglement

# Problem: Entanglement revivals

2 non-interacting qubits



Entanglement A-B



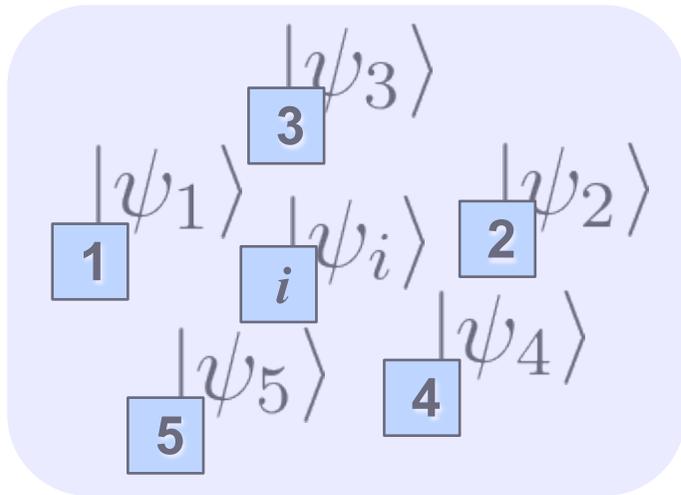
How does entanglement revives?

No non-local operations

No exchange of quantum correlations system-environment

# Entanglement of a mixed state

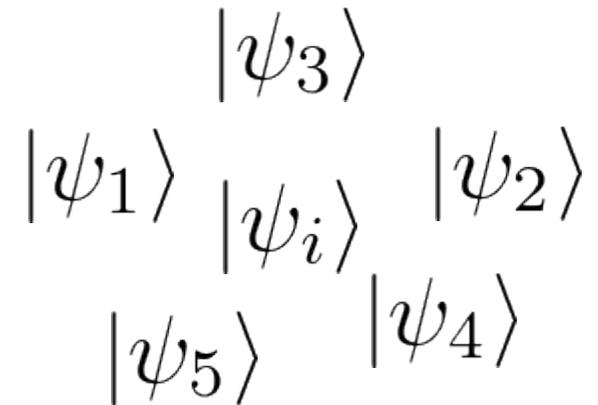
Charlie



Alice and Bob



Charlie sends the prepared systems to Alice and Bob



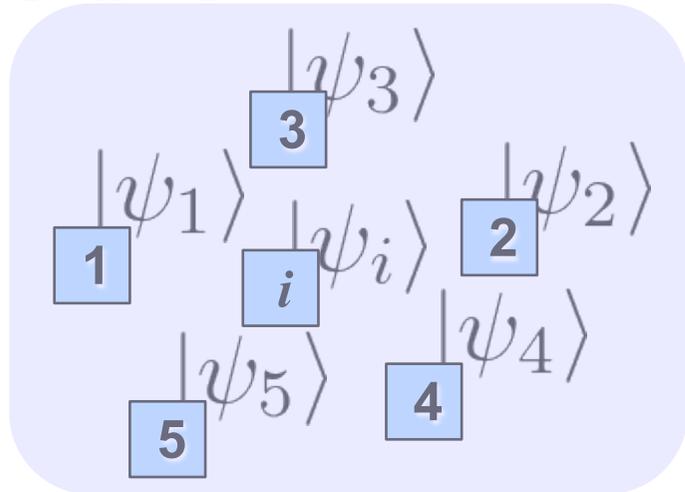
$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

How much is the entanglement shared by Alice and Bob?

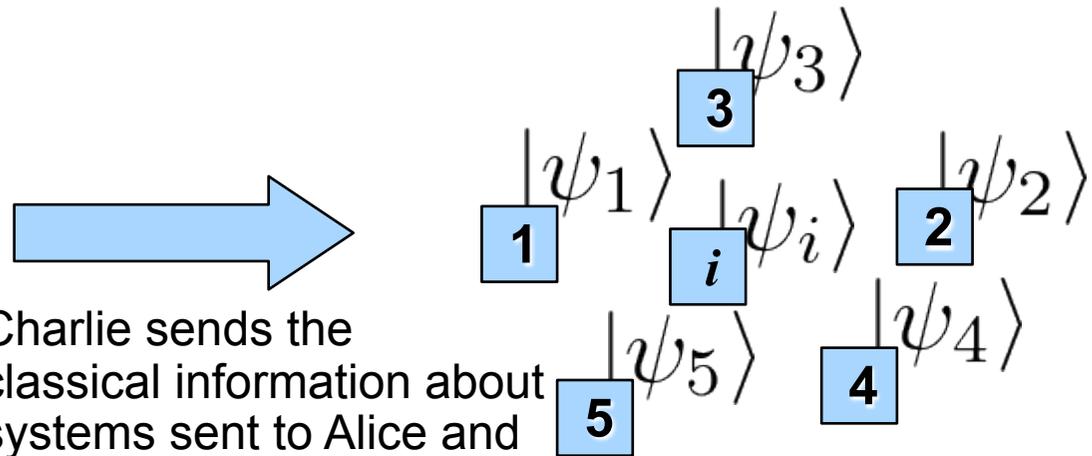
$$E(\rho)$$

# Entanglement of a quantum ensemble

Charlie



Alice and Bob



Charlie sends the classical information about systems sent to Alice and Bob

What changes if Alice and Bob try to distill entanglement?

$$E_{av}(\mathcal{A}) = \sum p_i E(|\psi_i\rangle\langle\psi_i|)$$

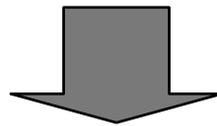
$$\mathcal{A} = \{(p_i, |\psi_i\rangle)\}$$

## Definition: Hidden entanglement

$$\begin{aligned} E_h(\mathcal{A}) &\equiv E_{av}(\mathcal{A}) - E(\rho) = \\ &= \sum_i p_i E(|\psi_i\rangle\langle\psi_i|) - E\left(\sum_i p_i |\psi_i\rangle\langle\psi_i|\right) \end{aligned}$$

$E$  is any convex entanglement measure, so that  $E_h \geq 0$

$E_h$ : **Entanglement** not exploitable, lack of classical information



It may be recovered **without non-local operations**

No complete knowledge...

$$\begin{array}{c} |\psi_3\rangle \\ |\psi_1\rangle \quad |\psi_i\rangle \quad |\psi_2\rangle \\ |\psi_5\rangle \quad |\psi_4\rangle \\ \rho = \sum_i p_i |\psi_i\rangle \langle \psi_i| \end{array}$$

classical  
information  
about systems

$$\begin{array}{c} |\psi_3\rangle \\ \boxed{3} \\ |\psi_1\rangle \quad |\psi_i\rangle \quad |\psi_2\rangle \\ \boxed{1} \quad \boxed{i} \quad \boxed{2} \\ |\psi_5\rangle \quad |\psi_4\rangle \\ \boxed{5} \quad \boxed{4} \\ \mathcal{A} = \{(p_i, |\psi_i\rangle)\} \end{array}$$

$$E(\rho) = E\left(\sum p_i |\psi_i\rangle \langle \psi_i|\right)$$

no non-local operations

no quantum correlations transfer

only classical information

$$E_{av}(\mathcal{A}) = \sum p_i E(|\psi_i\rangle \langle \psi_i|)$$

$$E_{av}(\mathcal{A}) - E(\rho) \geq 0$$

Entanglement increases

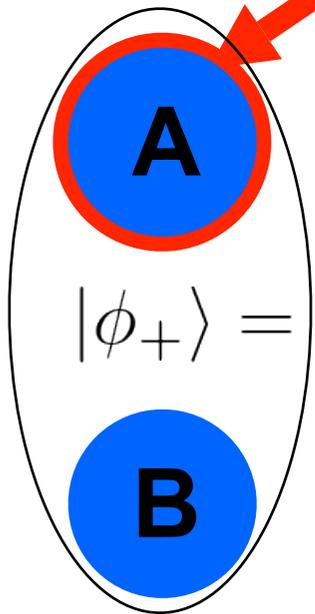
# Example: random local field

## Random Local Field

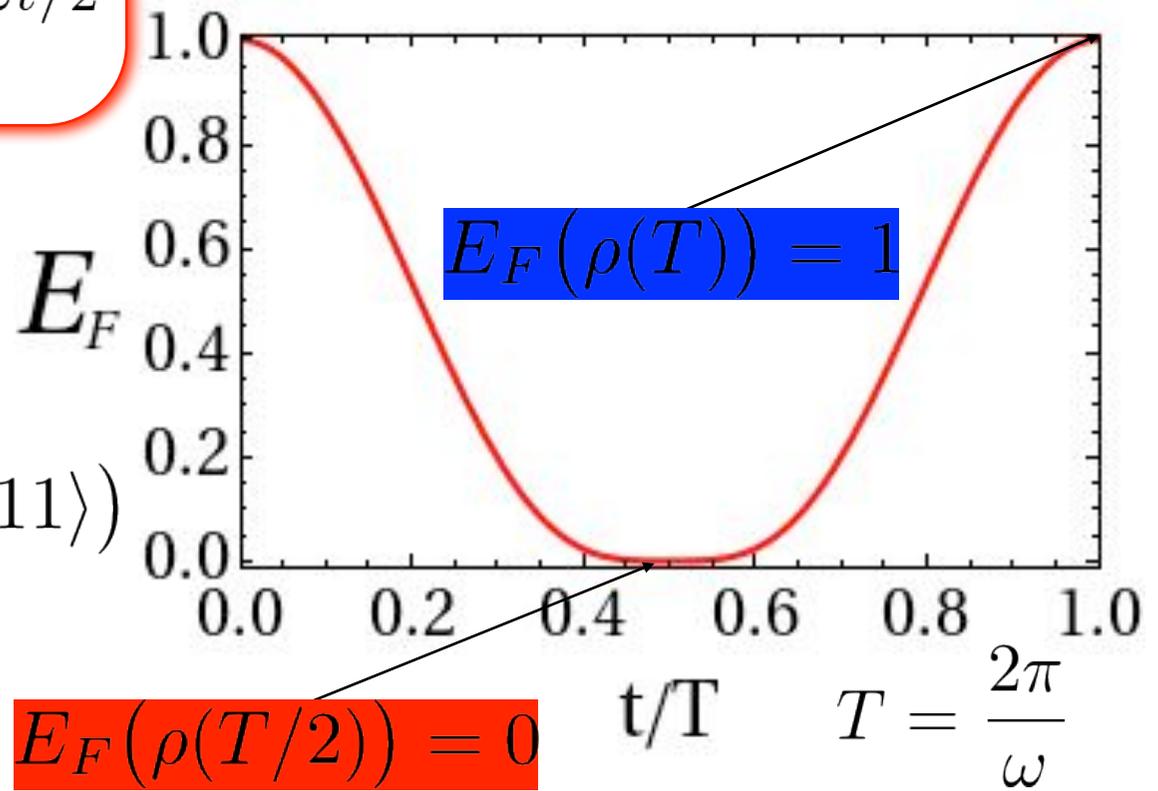
$$p_x = \frac{1}{2}, \quad U_x = e^{-i\sigma_x \omega t/2}$$

$$p_z = \frac{1}{2}, \quad U_z = e^{-i\sigma_z \omega t/2}$$

**entanglement increases**  
without **non-local operations**



$$|\phi_+\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$



$$E_F(\rho(T/2)) = 0 \xrightarrow{]0.5T, T]} E_F(\rho(T)) = 1$$

Manifestation of quantum correlations

that must be present before

## Quantum ensemble description

$$p_x, |\psi_x(t)\rangle = U_x|\phi_+\rangle$$

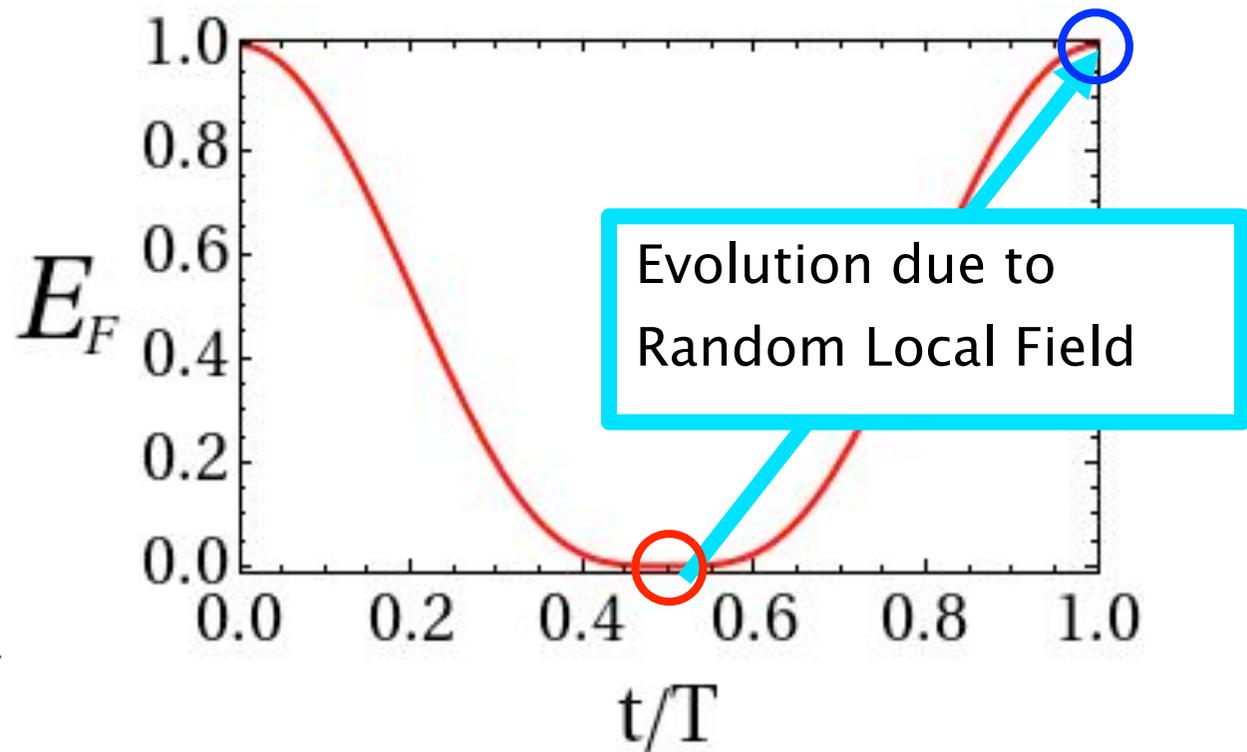
$$p_z, |\psi_z(t)\rangle = U_z|\phi_+\rangle$$

$$\mathcal{A}(t) = \{p_i, |\psi_i(t)\rangle\}$$

## Density operator formalism

$$\rho(t) = \sum_i p_i U_i |\phi_+\rangle \langle \phi_+| U_i^\dagger$$

mixture



$$E_{av}(\mathcal{A}) = \sum_i p_i E(|\psi_i(t)\rangle) = \sum_i p_i E(U_i|\phi_+\rangle) = 1 \quad \forall t$$

$$U_x(T/2) = -i\sigma_x \quad U_z(T/2) = -i\sigma_z$$

$$\mathcal{A}(T/2) = \left\{ \left( \frac{1}{2}, |\phi_-\rangle \right), \left( \frac{1}{2}, |\psi_+\rangle \right) \right\}$$

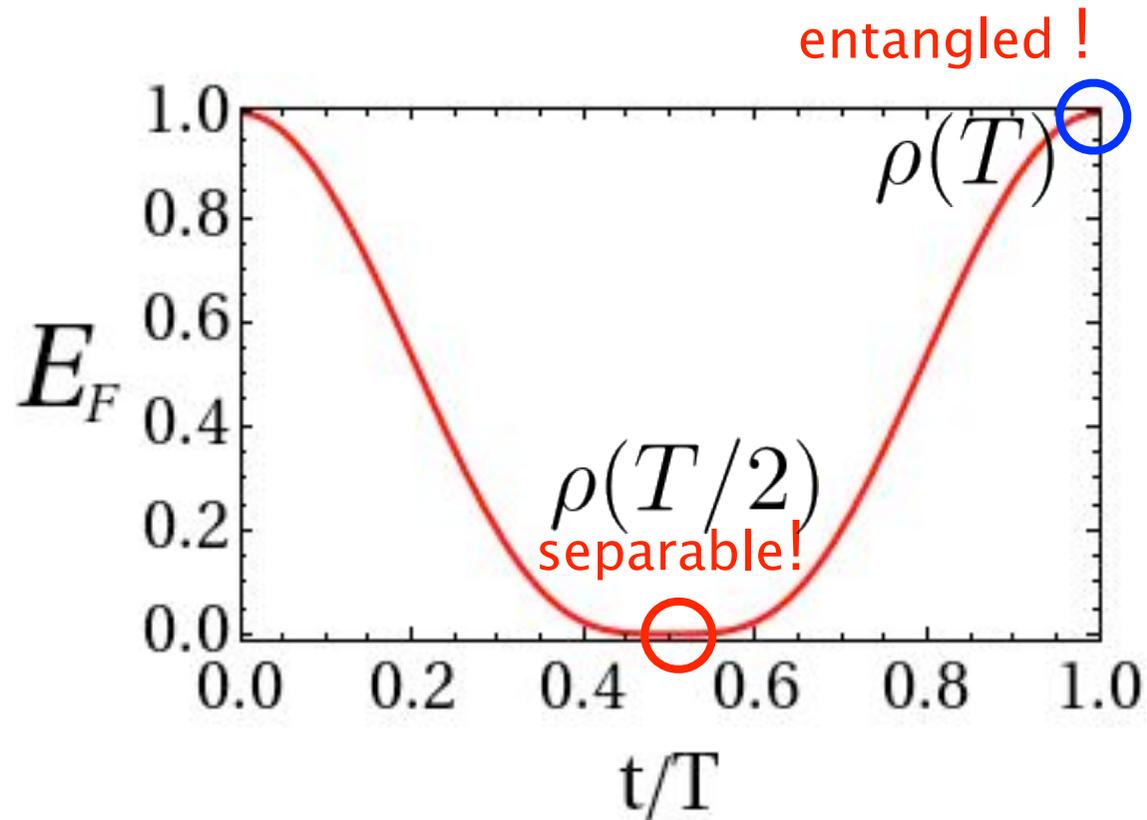
$$E(\rho(T/2)) = 0$$

$$E_h(\mathcal{A}(T/2)) = 1$$

$$U_x(T) = U_z(T) = I_d$$

$$\mathcal{A}(T) = \{1, |\phi_+\rangle\}$$

$$E(\rho(T)) = 1$$

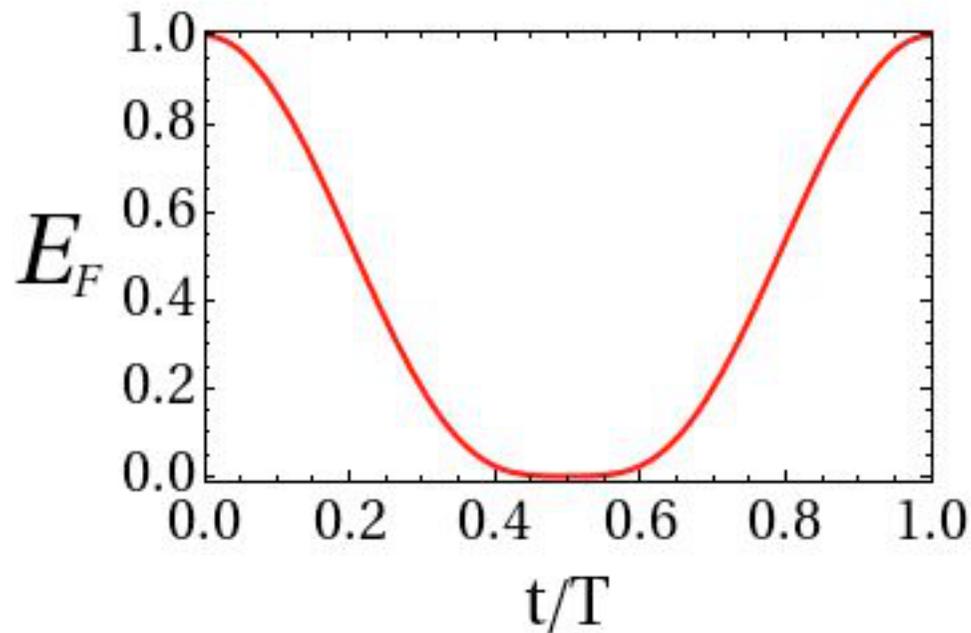


# Violation of the monotonicity axiom?

Monotonicity axiom:

*Entanglement cannot increase under LOCC*

The entanglement recovery is induced by local operations that are not LOCC (the evolution of the density operator from  $T/2$  to  $T$  is **not a CPT map**)



# Purely dephasing random telegraph noise

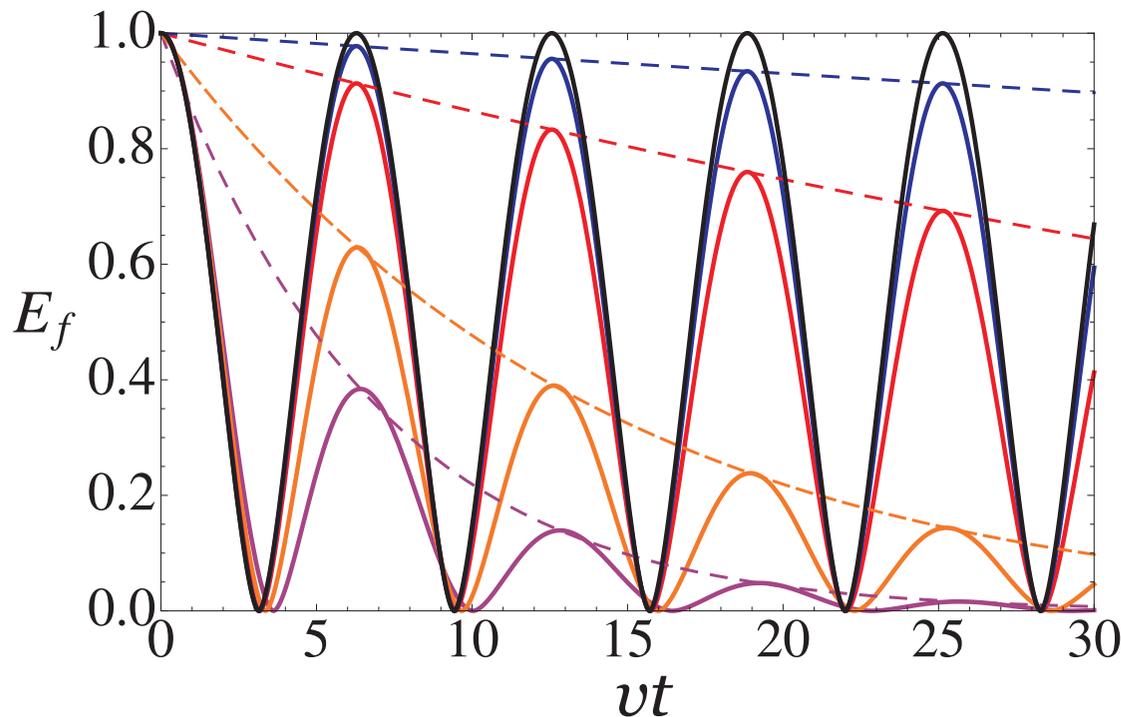
$$\mathcal{H} = \mathcal{H}_0 + \delta\mathcal{H},$$

$$\mathcal{H}_0 = -\frac{\Omega_A}{2}\sigma_{z_1} - \frac{\Omega_B}{2}\sigma_{z_2}, \quad \delta\mathcal{H} = -\frac{\xi(t)}{2}\sigma_{z_1} \quad \xi(t) \in \{0, v\}$$

switching rate  $\gamma$

At  $\gamma=0$ :  $\mathcal{A} = \left\{ \left( p_0, |\varphi_0(t)\rangle \right), \left( p_v, |\varphi_v(t)\rangle \right) \right\}$   $|\varphi_0(t)\rangle = \frac{1}{\sqrt{2}} \left( |00\rangle + e^{-i(\Omega_A+\Omega_B)t} |11\rangle \right)$

$$|\varphi_v(t)\rangle = \frac{1}{\sqrt{2}} \left( |00\rangle + e^{-ivt} e^{-i(\Omega_A+\Omega_B)t} |11\rangle \right)$$



Entanglement revivals at

$$t_n = 2n\pi/v$$

(for  $v/\gamma > 1$ )

Peak values drop as

$$f(C) = f(e^{-\gamma t/2})$$

# Stochastic low-frequency noise

## Stochastic low-frequency noise

$$\mathcal{H}_A = -\frac{\Omega_A}{2}\sigma_z - \frac{\varepsilon(t)}{2}\sigma_z$$

Static noise

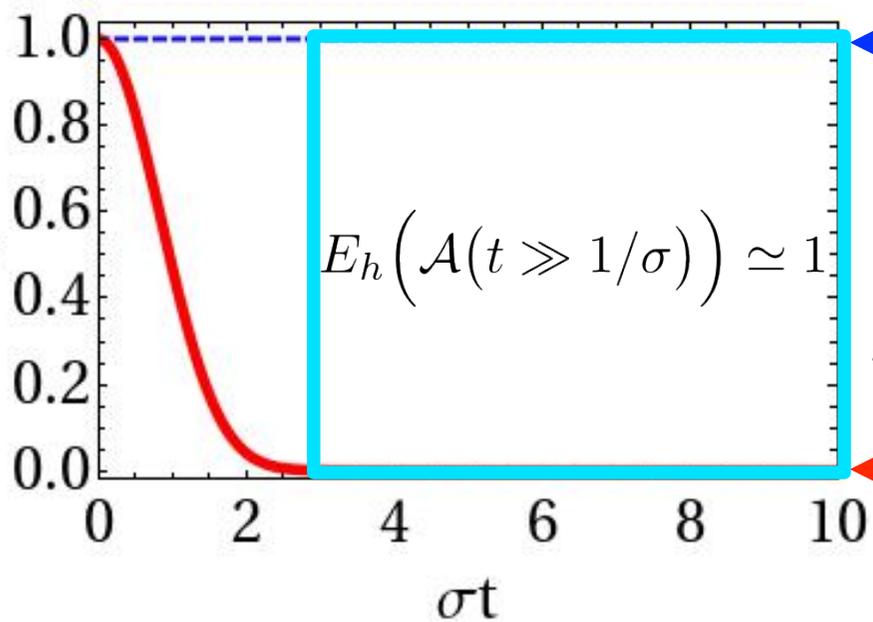
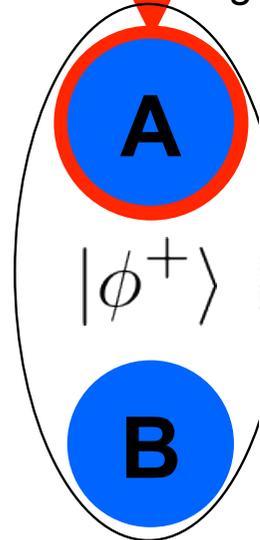
$$\varepsilon(t) \simeq \varepsilon \quad C(\rho(t)) = e^{-\frac{1}{2}\sigma^2 t^2}$$

gaussian, variance:  $\sigma^2$

For each realization of  $\varepsilon(t)$

$$|\psi_\varepsilon(t)\rangle = e^{-i\frac{\Omega_A + \varepsilon}{2}\sigma_z t} |\phi_+\rangle$$

$$\mathcal{A}(t) = \{p(\varepsilon)d\varepsilon, |\psi_\varepsilon(t)\rangle\}$$



$$E_{av}(\mathcal{A}(t)) = 1$$

$$\rho(t) = \int d\varepsilon p(\varepsilon) |\psi_\varepsilon(t)\rangle \langle \psi_\varepsilon(t)|$$

$$E_F(\rho(t \gg 1/\sigma)) \simeq 0$$

# Recovering entanglement by local pulses

## Stochastic low-frequency noise

$$\mathcal{H}_A = -\frac{\Omega_A}{2}\sigma_z - \frac{\varepsilon(t)}{2}\sigma_z + \mathcal{V}(t, \bar{t})\sigma_x$$

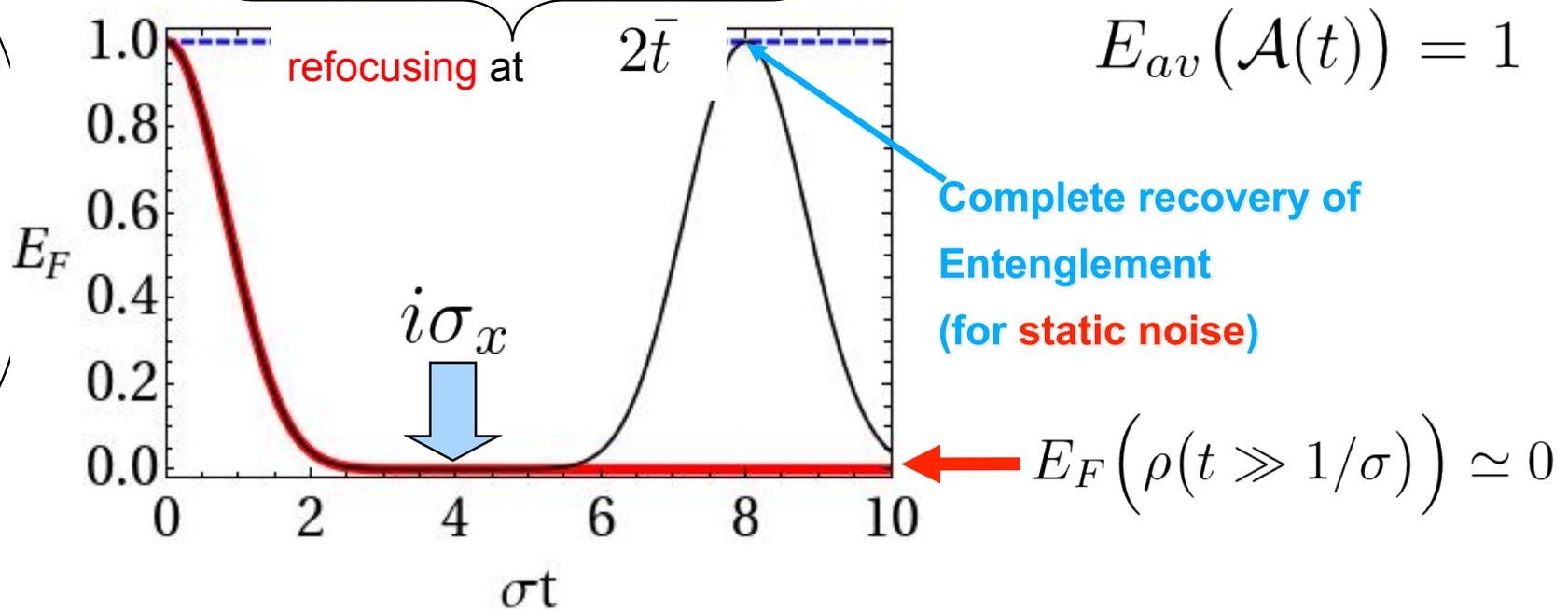
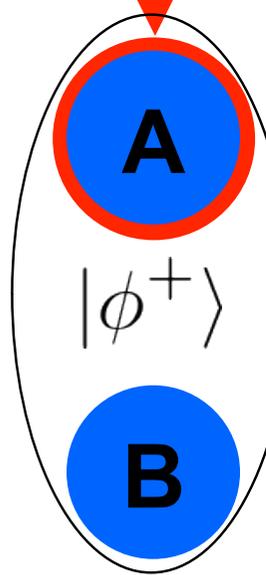
$$e^{-i\sigma_x\pi/2} = -i\sigma_x$$

$\pi$ -pulse  
around x

$$|\psi_\varepsilon(t)\rangle = e^{-i\frac{\Omega_A+\varepsilon}{2}\sigma_z(t-\bar{t})} \sigma_x e^{-i\frac{\Omega_A+\varepsilon}{2}\sigma_z\bar{t}} |\psi(0)\rangle$$

$$= e^{-i\frac{\Omega_A+\varepsilon}{2}\sigma_z(t-\bar{t})} e^{+i\frac{\Omega_A+\varepsilon}{2}\sigma_z\bar{t}} \sigma_x |\psi(0)\rangle$$

$i\sigma_x$

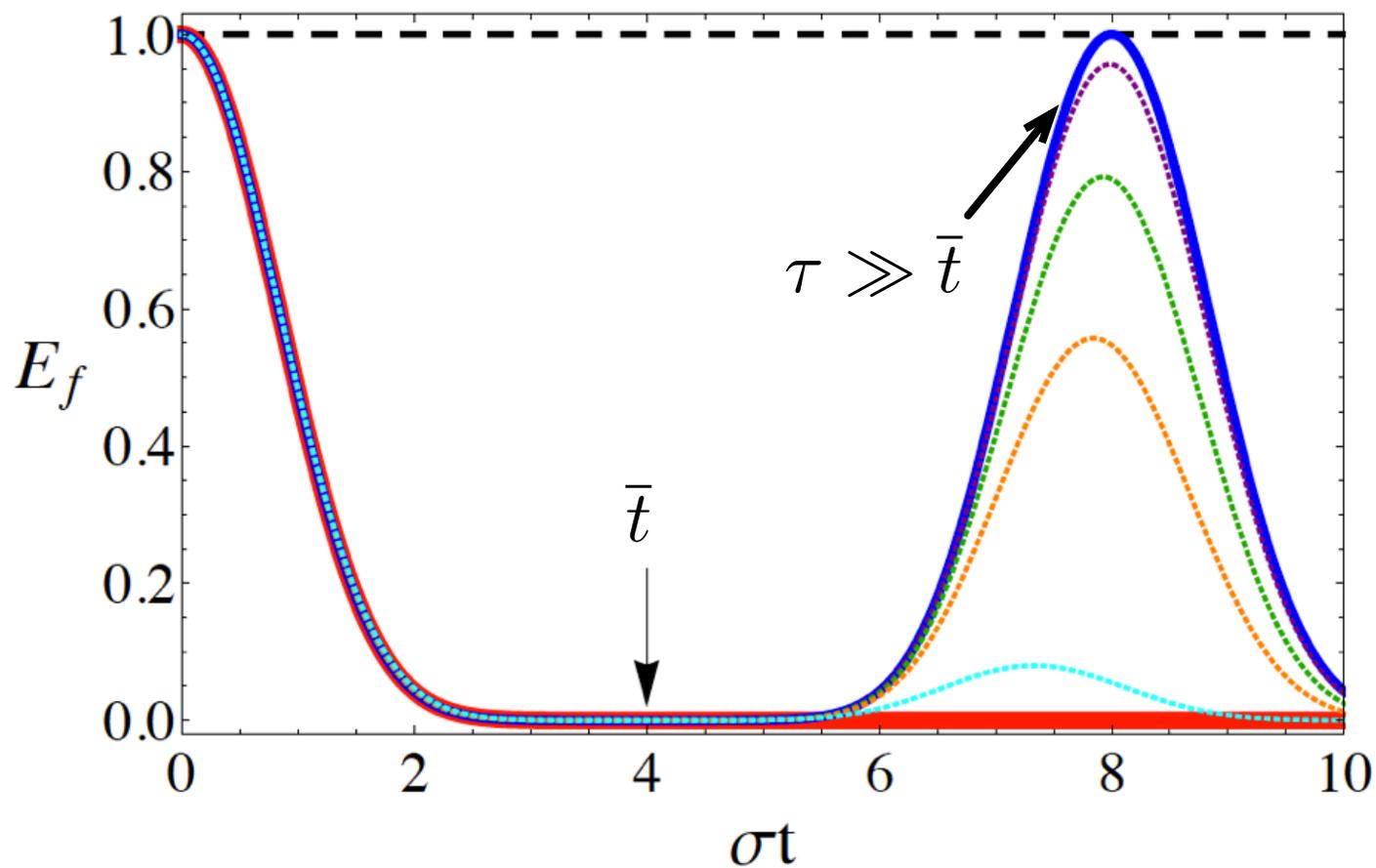


$$E_{av}(\mathcal{A}(t)) = 1$$

$$E_F(\rho(t \gg 1/\sigma)) \simeq 0$$

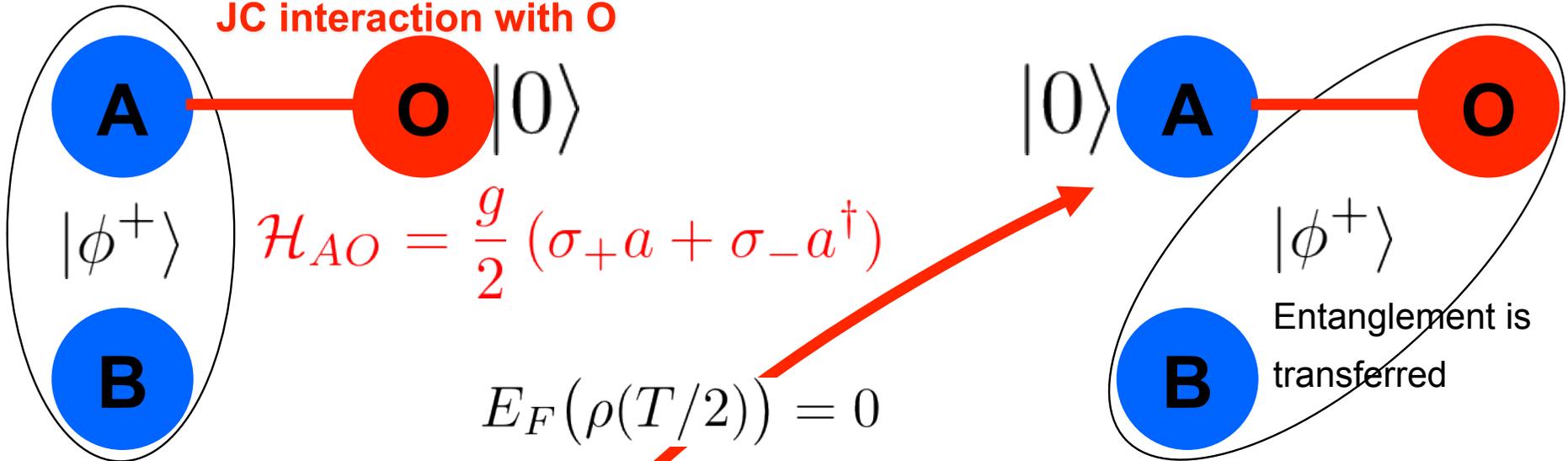
# Partial recovery with **finite noise correlation time scale**

$$\langle \varepsilon(t)\varepsilon(0) \rangle = \sigma^2 e^{-|t|/\tau}$$

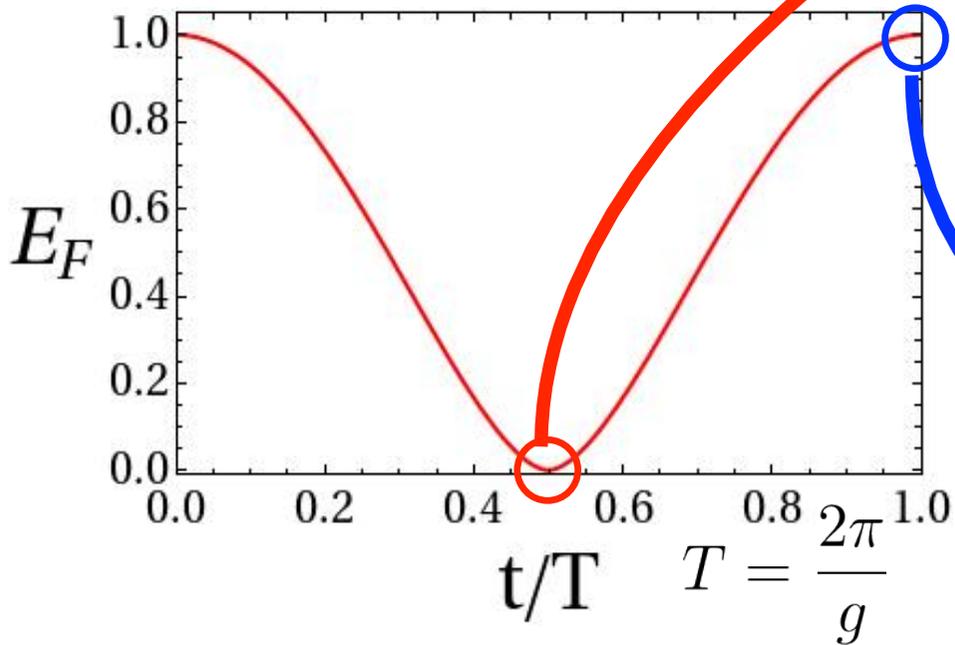


# Non-Markovian dynamics and entanglement revivals

JC interaction with O



$$E_F(\rho(T/2)) = 0$$



$$\mathcal{A}(T/2) = \{p_x, |0_A\rangle \otimes |x_B\rangle\} \quad E_{av} = 0$$

$$E_h = E_{av} - E_f = 0$$

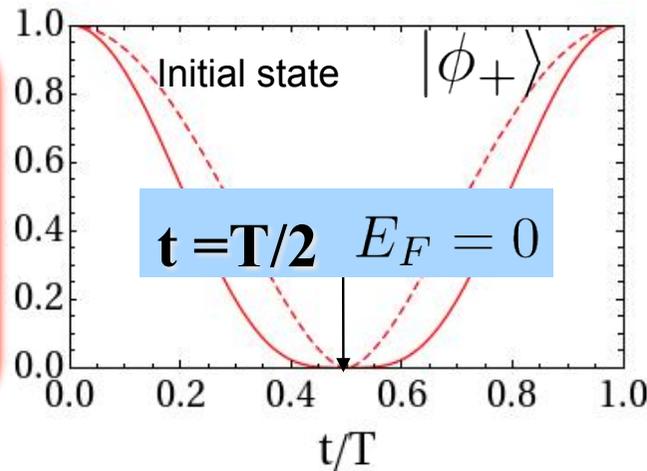
Perfect entanglement back-transfer  
from OB to AB

# Entanglement revivals in classical vs quantum environments

**Random Local Field**

$$p_z = \frac{1}{2}, \quad U_z = -i\sigma_z$$

$$p_x = \frac{1}{2}, \quad U_x = -i\sigma_x$$



**JC interaction with O**

$$\mathcal{H}_{AO} = \frac{g}{2} (\sigma_+ a + \sigma_- a^\dagger)$$

SWAP(A,O)

$$\mathcal{A}(T/2) = \left\{ \left( \frac{1}{2}, |\phi_-\rangle \right), \left( \frac{1}{2}, |\psi_+\rangle \right) \right\}$$

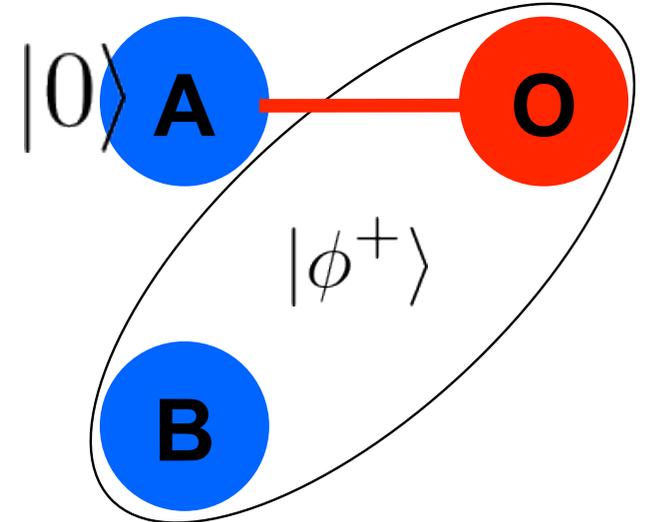
$$\mathcal{A}(T/2) = \{ p_x, |0_A\rangle \otimes |x_B\rangle \}$$

$$E_{av} = 1$$

$$E_{av} = 0$$

$$E_h = 1$$

$$E_h = 0$$



classical information about  
which random unitaries AB  
underwent

**Recovery!**

**No recovery!**

# HE in systems with damping

Assume that one qubit of a Bell pair is subject to damping (quantum environment) modeled by a Lindblad master equation

$$\dot{\rho} = -\frac{i}{\hbar} [H, \rho] - \frac{\Gamma}{2} (\sigma_- \sigma_+ \rho + \rho \sigma_- \sigma_+) + \Gamma \sigma_+ \rho \sigma_-$$

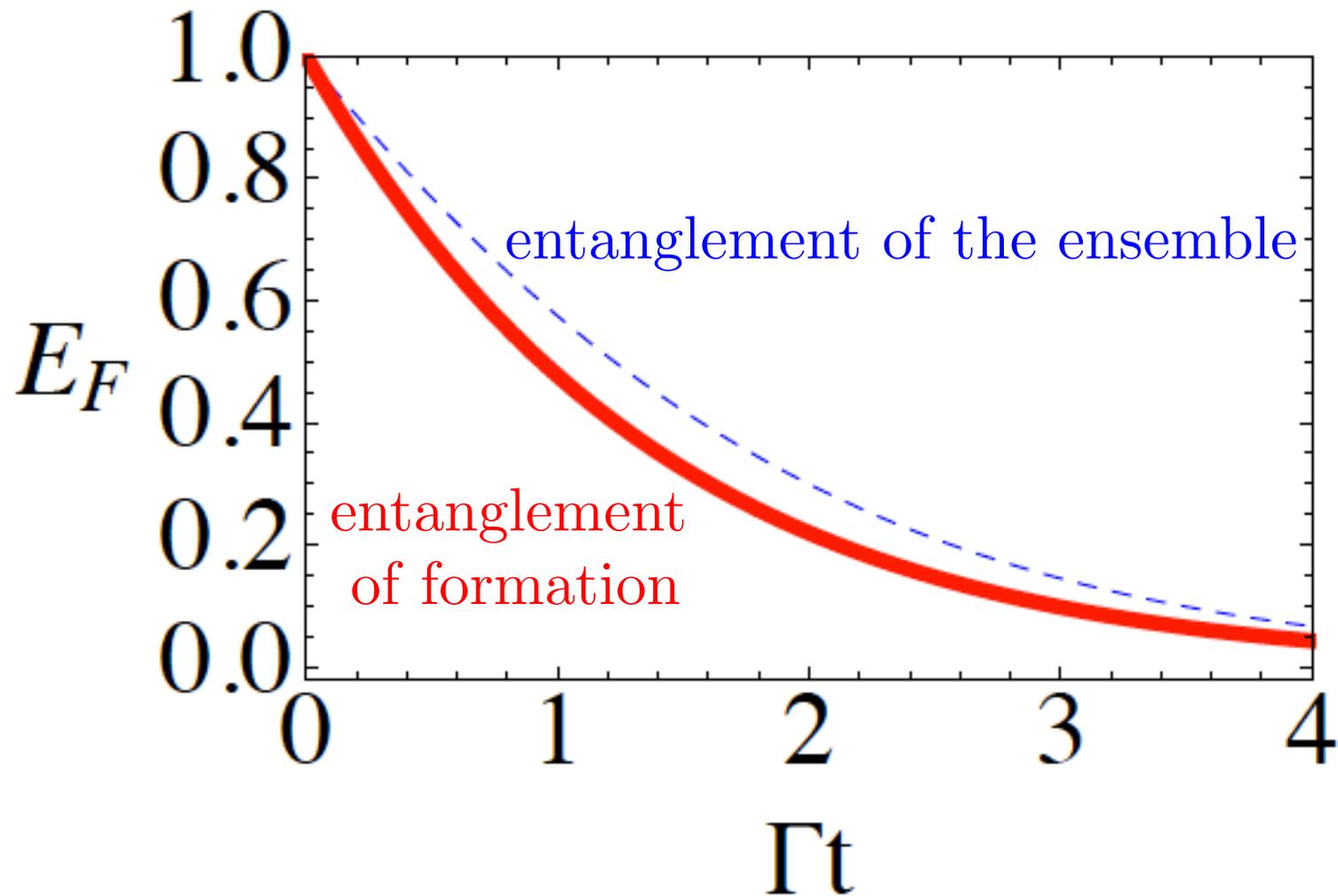
Relevant quantum ensemble when the emitted photons are detected

$$\mathcal{A}(t) = \{(p_0(t), |\varphi_0(t)\rangle), (p_1(t), |\varphi_1(t)\rangle)\}$$

$$|\varphi_0(t)\rangle = |10\rangle,$$

$$|\varphi_1(t)\rangle = \frac{1}{\sqrt{1 + e^{-\Gamma t}}} \left( e^{i\frac{\Omega_A + \Omega_B}{2} t} |00\rangle + e^{-i\frac{\Omega_A + \Omega_B}{2} t} e^{-\frac{\Gamma}{2} t} |11\rangle \right)$$

$$p_0(t) = \frac{1}{2}(1 - e^{-\Gamma t}), \quad p_1(t) = \frac{1}{2}(1 + e^{-\Gamma t})$$



To recover the hidden entanglement it is sufficient to discard the AB systems each time a photon is detected (**classical information**)

# Summary

Defined “**Hidden**” **Entanglement (HE)** on the basis of the ensemble description of the system dynamics

**HE**>0: a recovery of **entanglement** by local operation and classical communication is possible, even in realistic cases in which it is **not possible to control the environment** (HE conceptually different from the entanglement of assistance)

Practical relevance in solid state quantum computing: **HE** indicates the amount of entanglement recoverable by local pulses (**dynamical decoupling**: echo, bang-bang....)

**Non-Markovian** dynamics:  
classical vs quantum environments