

# “Transversity” Correlations in Azimuthal and Single Spin Asymmetries

Leonard Gamberg\*

Division of Science, Penn State Berks

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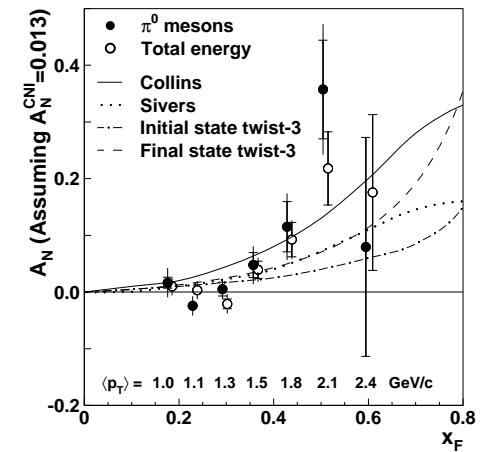
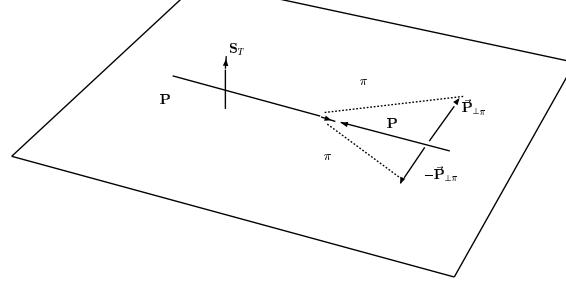
- Remarks on TSSA: QCD Correlations btwn. intrinsic  $k_{\perp}$  and transverse spin  $S_T$  properties of hadrons and quarks
- \* “Novel” Transversity Properties in Hard Scattering
- \* Reaction Mechanism-ISI/FSI: “ $T$ -odd” Structure and Fragmentation Functions and role in TSSA and AA
- \* Estimates of the Collins and Sivers Asymmetries
- \* Double  $T$ -odd  $\cos 2\phi$  asymmetry: SIDIS (DRELL-YAN-talk of G. Goldstein)
- \* Status: Investigation of Collins Function Spec. Model
- \* Other Method to Access Transversity
- Conclusions

\*G. R. Goldstein, Tufts Univ., K.A. Oganessyan *Financial District NYC*, and D.S. Hwang Sejong Univ, Korea  
Transversity 2005 COMO 9<sup>th</sup> Sept 2005

# Transverse SSA (TSSA) and AZIMUTHAL ASYMMETRIES (AA)

- **LARGE TSSAS OBSERVED**

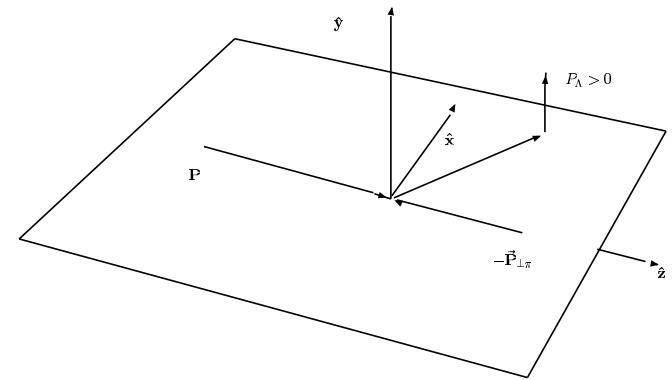
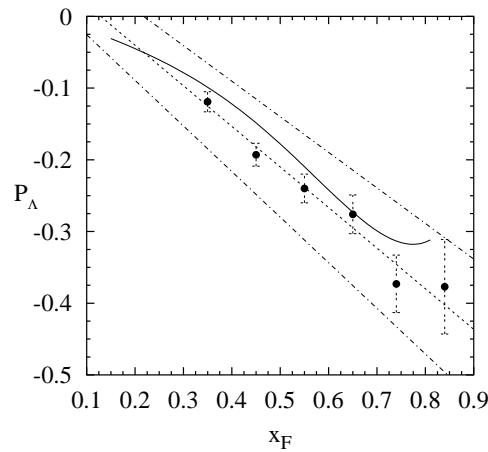
$$A_N = \frac{d\sigma^{p \uparrow p \rightarrow \pi} X - d\sigma^{p \downarrow p \rightarrow \pi} X}{d\sigma^{p \uparrow p \rightarrow \pi} X + d\sigma^{p \downarrow p \rightarrow \pi} X}$$



L-R asymmetry of  $\pi$  production and  $A_N$  for  $\pi^0$  production at STAR : PRL 2004

$P_\Lambda$  in p-p scattering from Fermi Lab

$$P_\Lambda = \frac{d\sigma^{pp \rightarrow \Lambda \uparrow X} - d\sigma^{pp \rightarrow \Lambda \downarrow X}}{d\sigma^{pp \rightarrow \Lambda \uparrow X} + d\sigma^{pp \rightarrow \Lambda \downarrow X}}$$



Heller,...,Bunce PRL:1983 PRL: 1983: Up-down asymmetry depicted for  $\Lambda$  production in p-p COM-frame.

## Transverse SSA (TSSA) and AZIMUTHAL ASYMMETRIES (AA)

★ Colinear approximation

TSSA vanishingly small at large scales and leading order  $\alpha_s$   
Generically,

$$|\perp/\tau\rangle = \frac{1}{\sqrt{2}}(|+\rangle \pm i|-\rangle) \Rightarrow A_N = \frac{d\hat{\sigma}^\perp - d\hat{\sigma}^\tau}{d\hat{\sigma}^\perp + d\hat{\sigma}^\tau} \sim \frac{2 \operatorname{Im} f^* f^-}{|f^+|^2 + |f^-|^2}$$

- ★ Requires *helicity flip* as well as relative phase btwn helicity amps
- At partonic level massless QCD conserves helicity Born amplitudes are real!
- ★ Interference btwn loops-tree level Kane, Repko, PRL:1978 yield  $A_N \sim m_q \alpha_s / \sqrt{s}$
- Twist three effects analyzed for moderately large  $p_T$

Efremov-Teryaev:PLB 1985, Qiu Sterman:PRL 1991, and Koike:PLB 2000

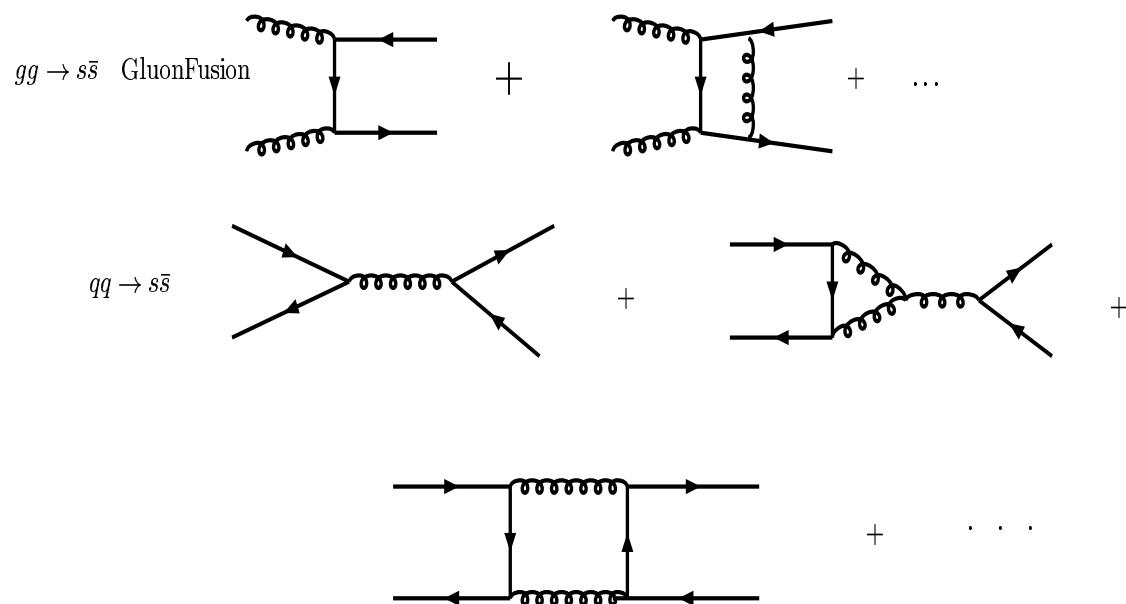
## Inclusive $\Lambda$ production From PQCD ( $pp \rightarrow \Lambda^\uparrow X$ )

$$P_\Lambda = \frac{d\sigma^{pp \rightarrow \Lambda^\uparrow X} - d\sigma^{pp \rightarrow \Lambda^\downarrow X}}{d\sigma^{pp \rightarrow \Lambda^\uparrow X} + d\sigma^{pp \rightarrow \Lambda^\downarrow X}}$$

- Need a strange quark to Polarize a  $\Lambda$  ( $pp \rightarrow \Lambda^\uparrow X$ ) PQCD contributions calculated

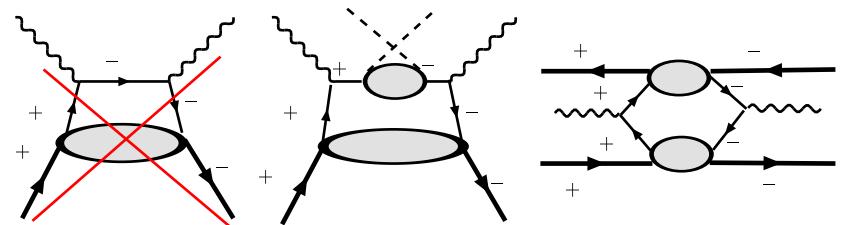
Dharmarajna & Goldstein PRD 1990

$P_\Lambda$  goes like  $m_q \alpha_s / \sqrt{s}$  as predicted  $m_q$  is the strange quark mass



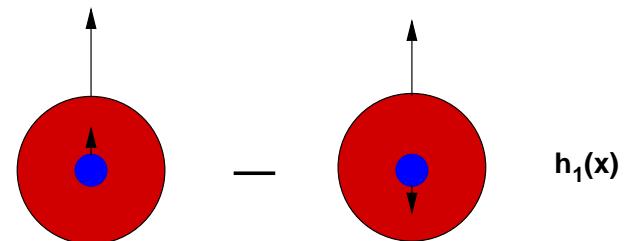
# Helicity Flips Accommodated in Hard Scattering, from “Transversity” Distributions

Drell-Yan  $p_\perp p_\perp \Rightarrow l^+ l^- X$  (2 in the initial)  
 SIDIS  $l p_\perp \Rightarrow l' h X$  (1 in initial 1 in final)



\* DY: Ralston and Soper NPB:1979 encountered double transverse spin asymmetry

$$A_{TT}^{DY} = \frac{2 \sin^2 \theta \cos(\phi_1 + \phi_2)}{1 + \cos^2 \theta} \frac{\sum_a e_a^2 h_1^a(x) \bar{h}_1^a(x)}{\sum_a e_a^2 f_1^a(x) \bar{f}_1^a(x)}$$



$h_1(x)$  probability to find quark with spin polarized along transverse spin direction minus oppositely polarized case

# Transversity and Twist-3 Contributions to Asymmetries

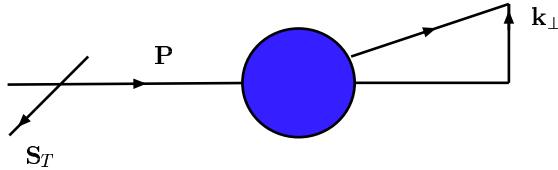
- ★ SIDIS: Jaffe and Ji PRL:1993 encountered at twist three level double spin asymmetry Estimate, Gamberg, Hwang, Oganessyan PLB:2004

$$A_{LT} = \frac{\lambda_e |S_T| \sqrt{1-y} \frac{4}{Q} [M x g_T(x) D_1(z) + M_h h_1(x) \frac{E(z)}{z}]}{\frac{[1+(1-y)^2]}{y} f_1(x) D_1(z)}$$

- ★ Analogous process in Drell-Yan  $\pi P \rightarrow \mu^+ \mu^- X$   
Ji PLB:1992

## “ $T$ -Odd” (or $A_\tau$ ) Correlations: Beyond Co-linearity

- TSSA indicative “ $T$ -odd” correlations among *transverse* spin and momenta  
e.g.  $P P^\perp \rightarrow \pi X \quad S_T \cdot (P \times k_\perp)$

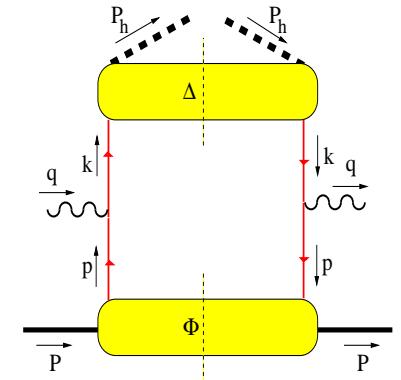


- Sensitivity to  $k_\perp$  intrinsic quark momenta, associated non-perturbative transverse momentum distribution functions **TMD**  
Soper, PRL:1979:  $\int d\mathbf{k}_\perp \mathcal{P}(\mathbf{k}_\perp, x) = f(x)$
- Correlation accounts for left-right transverse SSA Sivers: PRD 1990 in inclusive  $\pi$  production (Anselmino & Murgia PLB: 1995 ...)
- Collins NPB 1993 proposed  $T$ -odd correlation of transversely polarized fragmenting quark: TSSA in lepto-production  $\ell \vec{p} \rightarrow \ell' \pi X$   
Initial-Final state effect:  $s_T \cdot (p \times P_{h\perp})$ ,  $s_T$  is the spin of fragmenting quark,  $p$  is quark momentum and  $P_{h\perp}$  is transverse momentum produced pion

## T-Odd Correlations: Beyond Co-linearity

Recent Times Boer & Mulders and Co. incorporated  $k_{\perp}$  T-odd PDFs and FFs:  
 Relevant to hard scattering QCD at leading twist. Adopted Factorized Description  
 Ellis, Furmanski, Petronzio NPB: 1982, Collins *et al.* PQCD... : 82 , J. Qui PRD: 1990, Levelt & Mulders, Mulders &  
 Tangerman, NPB: 1994, 1996

$$\frac{d\sigma^{\ell N \rightarrow \ell' h X}}{dxdydz d^2 P_{h\perp}} = \frac{M\pi\alpha^2 y}{2Q^4 z} L_{\mu\nu} \mathcal{W}^{\mu\nu}$$



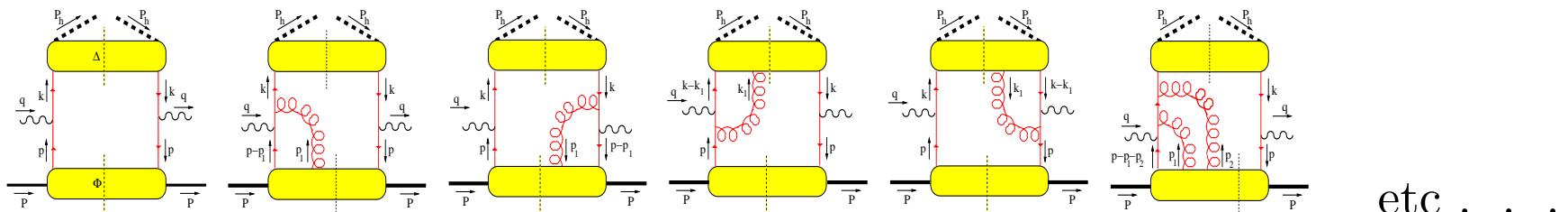
### Hadronic Tensor

$$2M\mathcal{W}^{\mu\nu}(q, P, P_h) = \int d^2 \mathbf{p}_T d^2 \mathbf{k}_T \delta^2(\mathbf{p}_T + \mathbf{q}_T - \mathbf{k}_T) \text{Tr}[\Phi(x_B, \mathbf{p}_T) \gamma^\mu \Delta(z_h, \mathbf{k}_T) \gamma^\nu] \\ + (q \leftrightarrow -q, \mu \leftrightarrow \nu)$$

# Color Gauge Invariance Built into Factorized QCD at “leading twist”-Wilson Line & T-Odd Contributions to QCD Processes

- Gauge Invariant Distribution and Fragmentation Functions

Boer, Mulder: NPB 2000, Ji, Yuan & Belitsky PLB: 2002, NPB 2003, Boer, Mulder, Pijlman NPB 2003



*sub-class of loops* in eikonal limit (soft gluons) sum up to yield color gauge invariant hadronic tensor factorized into the distribution and fragmentation operators

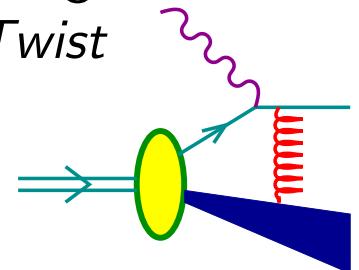
$$\Phi(p, P) = \int \frac{d^3 \xi}{2(2\pi)^3} e^{ip \cdot \xi} \langle P | \bar{\psi}(\xi^-, \xi_\perp) \mathcal{G}_{[\xi^-, \infty]}^\dagger | X \rangle \langle X | \mathcal{G}_{[0, \infty]} \psi(0) | P \rangle |_{\xi^+ = 0}$$

$$\Delta(k, P_h) = \int \frac{d^3 \xi}{4z(2\pi)^3} e^{ik \cdot \xi} \langle 0 | \mathcal{G}_{[\xi^+, -\infty]} \psi(\xi) | X; P_h \rangle \langle X; P_h | \bar{\psi}(0) \mathcal{G}_{[0, -\infty]}^\dagger | 0 \rangle |_{\xi^- = 0}$$

$$\mathcal{G}_{[\xi, \infty]} = \mathcal{G}_{[\xi_T, \infty]} \mathcal{G}_{[\xi^-, \infty]}, \quad \text{where} \quad \mathcal{G}_{[\xi^-, \infty]} = \mathcal{P} \exp(-ig \int_{\xi^-}^{\infty} d\xi^- A^+)$$

# Rescattering-ISI/FSI $T$ -Odd Contributions to Asymmetries

PLB: 2002 Brodsky, Hwang, and Schmidt demonstrate rescattering of a gluon could produce the necessary phase leading to nonzero SSAs at *Leading Twist*



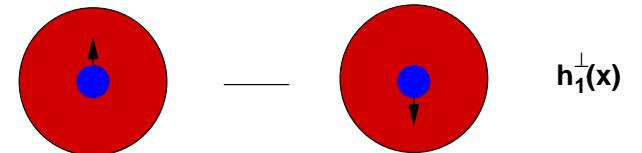
Initial-Final state effect:  $\mathbf{S}_T \cdot (\mathbf{P} \times \mathbf{k}_\perp)$

- Ji, Yuan PLB: 2002 describe effect in terms of gauge invariant distribution functions
- Demonstrates that BHS calculated Sivers Function  $f_{1T}^\perp(x, k_\perp)|_{\text{SDIS}}$   
In Singular gauge,  $A^+ = 0$ , **effect remains**
- Collins, PLB: 2002, modifies earlier claim of trivial Sivers Effect  
 $f_{1T}^\perp(x, k_\perp)|_{\text{SDIS}} = -f_{1T}^\perp(x, k_\perp)|_{\text{DY}}$

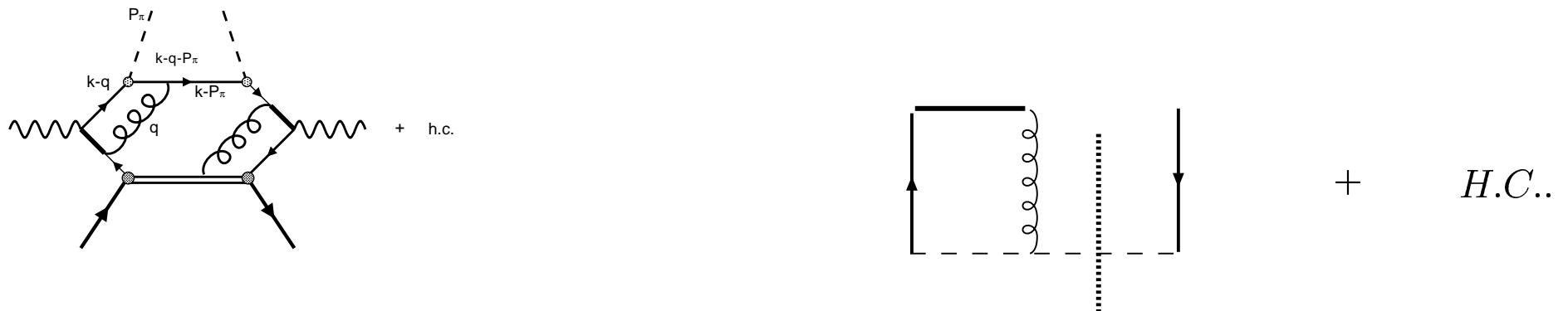
# FSI Mechanism can Generate Boer-Mulders- $h_1^\perp$

Goldstein, Gamberg–ICHEP-proc., Amsterdam: 2002, hep-ph/0209085, G, G and Oganessyan PRD 2003

- $h_1^\perp$  Naturally defined from gauge invariant TMD: Co-joined with  $H_1^\perp$  enters  $\cos 2\phi$  AA
- Applied “eikonal Feynman rules”  
to calculate (Collins, Soper, NPB: 1982)



$$\Phi_{[h_1^\perp]}^{[\sigma^{\perp+}\gamma_5]}(x, k_\perp) = \frac{1}{2} \int dp^- \text{Tr} \left( i\sigma^{+\perp}\gamma_5 \Phi \right) = \frac{\varepsilon_{+-\perp j} k_{\perp j}}{M} h_1^\perp(x, k_\perp)$$



$$\Phi^{[\Gamma]}(x, k_\perp) = \sum_X \int \frac{d\xi^- d^2\xi_\perp}{2(2\pi)^3} e^{-i\xi \cdot \vec{k}_\perp} \langle P | \bar{\psi}(\xi) \mathcal{G}_{[\xi^-, \infty]}^\dagger | X \rangle \langle X | \mathcal{G}_{[0, \infty]} \Gamma \psi(0) | P \rangle|_{\xi^+=0} + \text{h.c.}$$

$h_1^\perp(x, k_\perp)$ , represents, number density transversely polarized quarks in an unpolarized nucleons nucleons-complementary to  $f_{1T}^\perp(x, k_\perp)$ ,  
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# Provide source of T-Odd Contributions to TSSA and AA

- Enter the *leading twist distribution and fragmentation correlators “T-odd” Distribution Functions: Transversity Properties of quarks in Hadrons*

Boer, Mulder: PRD 1998

$$\Delta(z, \mathbf{k}_\perp) = \frac{1}{4} \left\{ D_1(z, z\mathbf{k}_\perp) \not{n}_- + H_1^\perp(z, z\mathbf{k}_\perp) \frac{\sigma^{\alpha\beta} k_{\perp\alpha} n_{-\beta}}{M_h} + D_{1T}^\perp \frac{\epsilon_{\mu\nu\rho\sigma} \gamma^\mu n_-^\nu k_\perp^\rho S_{hT}^\sigma}{M_h} + \dots \right\},$$

$$\Phi(x, \mathbf{p}_\perp) = \frac{1}{2} \left\{ f_1(x, \mathbf{p}_\perp) \not{n}_+ + h_1^\perp(x, \mathbf{p}_\perp) \frac{\sigma^{\alpha\beta} p_{T\alpha} n_{+\beta}}{M} + f_{1T}^\perp(x, \mathbf{p}_\perp) \frac{\epsilon^{\mu\nu\rho\sigma} \gamma^\mu n_+^\nu p_\perp^\rho S_T^\sigma}{M} + \dots \right\}$$

SIDIS cross section

$$d\sigma_{\{\lambda, \Lambda\}}^{\ell N \rightarrow \ell \pi X} \propto f_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 + \frac{k_\perp}{Q} f_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 \cdot \cos \phi$$

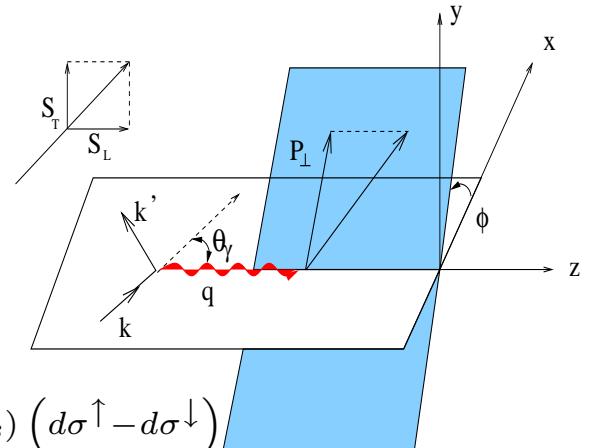
$$+ \left[ \frac{k_\perp^2}{Q^2} f_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 + h_1^\perp \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes H_1^\perp \right] \cdot \cos 2\phi$$

$$+ |S_T| \cdot h_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes H_1^\perp \cdot \sin(\phi + \phi_S) \quad \text{Collins}$$

$$+ |S_T| \cdot f_{1T}^\perp \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 \cdot \sin(\phi - \phi_S) \quad \text{Sivers}$$

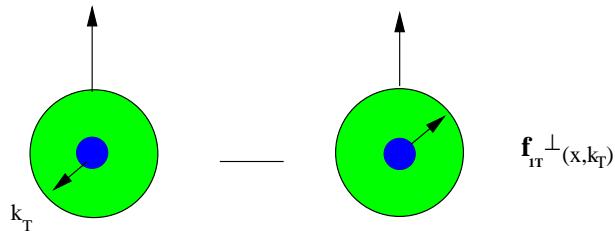
$$+ \dots$$

- Collins NPB:1993,  
Kotzinian NPB:1995, Mulders, Tangerman PLB:1995



$$\begin{aligned} \left\langle \frac{P_{h\perp}}{M\pi} \sin(\phi + \phi_s) \right\rangle_{UT} &= \frac{\int d\phi_s \int d^2 P_{h\perp} \frac{P_{h\perp}}{M\pi} \sin(\phi + \phi_s) (d\sigma^\uparrow - d\sigma^\downarrow)}{\int d\phi_s \int d^2 P_{h\perp} (d\sigma^\uparrow + d\sigma^\downarrow)} \\ &= |S_T| \frac{2(1-y) \sum_q e_q^2 h_1(x) z H_1^{\perp(1)}(z)}{(1+(1-y)^2) \sum_q e_q^2 f_1(x) D_1(z)} \end{aligned}$$

(Sivers PRD: 1990, Anselmino & Murgia PLB: 1995 ...)

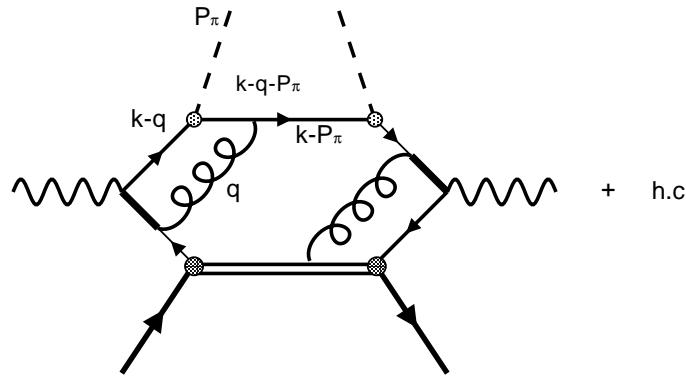


$$\left\langle \frac{|P_{h\perp}|}{M} \sin(\phi - \phi_S) \right\rangle_{UT} = |S_T| \frac{(1+(1-y)^2) \sum_q e_q^2 f_{1T}^{\perp(1)}(x) z D_1^q(z)}{(1+(1-y)^2) \sum_q e_q^2 f_1(x) D_1(z)},$$

- Probes the probability for a transversely polarized target, pions are produced asymmetrically about the transverse spin vector:

# $\cos 2\phi$ Asymmetry Generated by ISI & FSI thru Gauge link

Goldstein, Gamberg–ICHEP-Amsterdam: 2002, hep-ph/0209085, G,G, & Oganessyan PRD:2003



$$\begin{aligned}
 A_{UU}^{\cos(2\phi)} &= \left\langle \frac{|P_{h\perp}^2|}{MM_h} \cos 2\phi \right\rangle_{UU} \\
 &= \frac{8(1-y) \sum_q e_q^2 h_1^{\perp(1)}(x, Q^2) z^2 H_1^{\perp(1)q}(z, Q^2)}{(1+(1-y)^2) \sum_q e_q^2 f_1^q(x, Q^2) D_1^q(z, Q^2)}
 \end{aligned}$$

$$\frac{d\sigma}{dxdydzd^2P_\perp} \propto f_1 \otimes D_1 + \frac{k_T}{Q} f_1 \otimes D_1 \cdot \cos \phi + \left[ \frac{k_T^2}{Q^2} f_1 \otimes D_1 + h_1^\perp \otimes H_1^\perp \right] \cdot \cos 2\phi$$

D. Boer, P. Mulders, PRD: 1998

# Estimates of T-odd Contribution in SIDIS (& and Azimuthal Asymmetries Drell Yan (GSI program)

## $\cos 2\phi$ Asymmetry

- ★ The spectator model used in previous rescattering calculations assumes point-like nucleon-quark-diquark vertex, **leads to logarithmically divergent, asymmetries**

Goldstein, Gamberg, ICHEP 2002; hep-ph/0209085,

Gamberg, Goldstein, Oganessyan PRD 2003; Boer, Brodsky, Hwang, PRD: 2003(Drell-Yan)

$$\begin{aligned} h_1^\perp(x, k_\perp) &= f_{1T}^\perp(x, k_\perp) \\ &= \frac{g^2 e_1 e_2}{4\pi(2\pi)^3} \frac{(1-x)(m+xM)}{\Lambda(k_\perp^2)} \frac{M}{k_\perp^2} \ln \frac{\Lambda(k_\perp^2)}{\Lambda(0)} \end{aligned}$$

$$\Lambda(k_\perp^2) = k_\perp^2 + x(1-x) \left( -M^2 + \frac{m^2}{x} + \frac{\lambda^2}{1-x} \right)$$

- Asymmetry involves weighted function

$$h_1^{(1)\perp}(x) \equiv \int d^2 k_\perp \frac{k_\perp^2}{2M^2} h_1^\perp(x, k_\perp^2) \quad \text{diverges}$$

## Gaussian Distribution in $k_\perp$

**Log** divergence addressed by approximating the transverse momentum dependence of the quark-nucleon-vertex by a Gaussian distribution in  $k_\perp^2$ ,

Gamberg, Goldstein, Oganessyan, PRD 67 (2003)

$$\langle n | \psi(0) | P \rangle = \left( \frac{i}{\not{P} - m} \right) \Upsilon(k_\perp^2) U(P, S), \quad b \equiv \frac{1}{\langle k_\perp^2 \rangle}$$

where  $\Upsilon(k_\perp^2) = \mathcal{N} e^{-bk_\perp^2}$ .

$U(P, S)$  nucleon spinor, and quark propagator comes from untruncated quark line

$$h_1^\perp(x, k_\perp) = \frac{e_1 e_2 g^2}{2(2\pi)^4 \pi^2} \frac{b^2}{\Lambda(k_\perp^2)} \frac{(m + xM)(1 - x)}{k_\perp^2} \mathcal{R}(k_\perp^2, x) \quad (1)$$

with

$$\mathcal{R}(k_\perp^2, x) = \exp^{-2b(k_\perp^2 - \Lambda(0))} \left( \Gamma(0, 2b\Lambda(0)) - \Gamma(0, 2b\Lambda(k_\perp^2)) \right)$$

- $\lim \langle k_\perp^2 \rangle \rightarrow \infty$  width goes to infinity, regain *log* result

## INPUTS: Boer-Mulders and Unpolarized Structure Function $f_1(x)$

$$f_1(x) = \frac{g^2}{(2\pi)^2} \frac{b^2}{\pi^2} (1 - x) \cdot \left\{ \frac{(m + xM)^2 - \Lambda(0)}{\Lambda(0)} - \left[ 2b \left( (m + xM)^2 - \Lambda(0) \right) - 1 \right] e^{2b\Lambda(0)} \Gamma(0, 2b\Lambda(0)) \right\}$$

\* Normalization,  $\int_0^1 u(x) = 2$

$$\int_0^1 d(x) = 1$$

● Black curve-  $xu(x)$

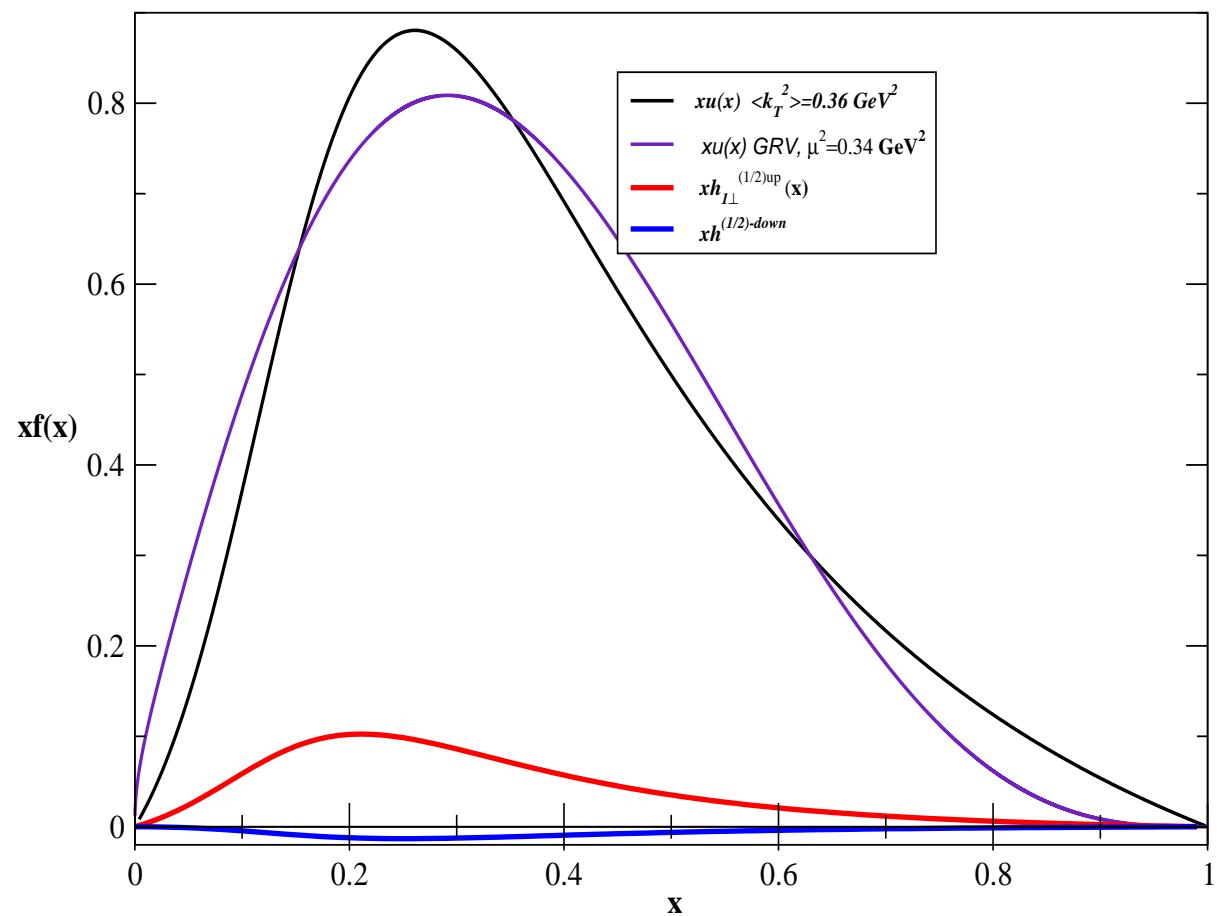
● Purple curve -  $xu(x)$  GRV

● Red curve  $xh_1^{\perp(1/2)(u)}$

● axial vector diquark coupling

Jakob, Mulders, Rodrigues NPB:199

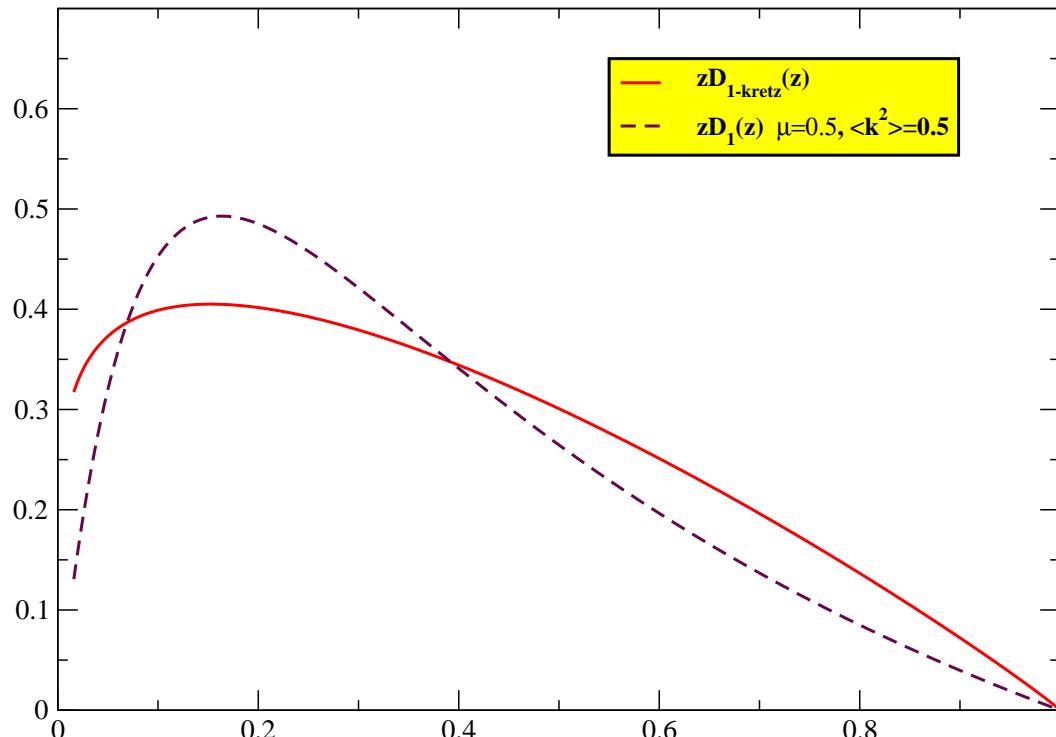
$$\gamma_5(\gamma^\mu + P^\mu/M)$$



## Pion Fragmentation Function

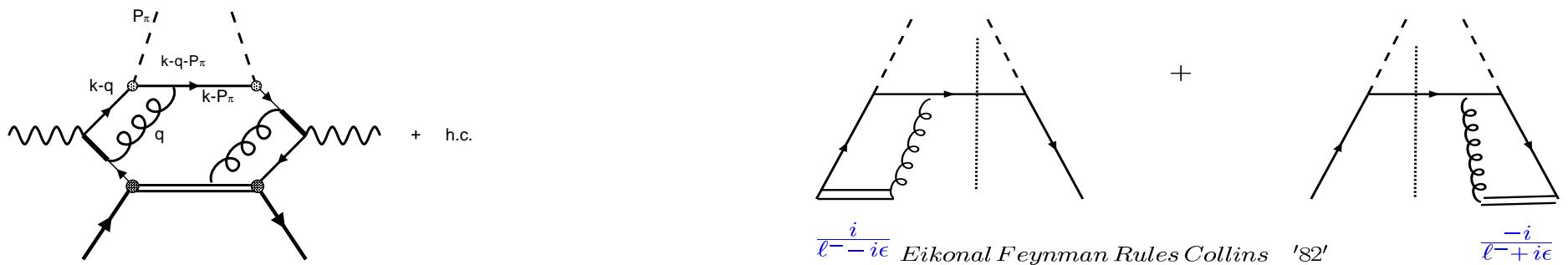
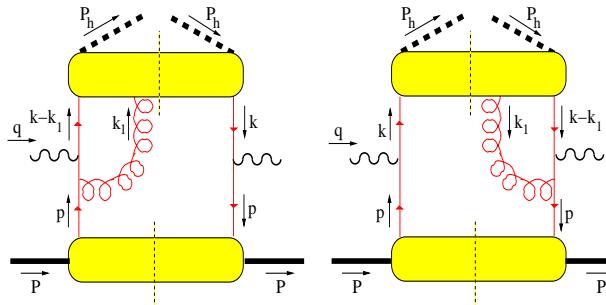
$$D_1(z) = \frac{N'^2 f_{qq\pi}^2}{4(2\pi)^2} \frac{1}{z} \frac{(1-z)}{z} \left\{ \frac{m^2 - \Lambda'(0)}{\Lambda'(0)} - \left[ 2b' (m^2 - \Lambda'(0)) - 1 \right] e^{2b' \Lambda'(0)} \Gamma(0, 2b' \Lambda'(0)) \right\},$$

which, multiplied by  $z$  at  $\langle k_\perp^2 \rangle = (0.5)^2 \text{ GeV}^2$  and  $\mu = m$ , estimates the distribution of Kretzer, PRD: 2000



# Gauge Link-Pole Contribution to T-Odd Collins Function

Gamberg,Goldstein,Oganessyan PRD68,2003  $\Delta^{[\sigma^\perp - \gamma_5]}(z, k_\perp) = \frac{1}{4z} \int dk^+ Tr(\gamma^- \gamma^\perp \gamma_5 \Delta)|_{k^- = P_\pi^- / z}$



Motivation: color gauge .inv frag. correlator “pole contribution”

We evaluate the projection  $\Delta^{[i\sigma^\perp - \gamma_5]}$ , results in leading twist, contribution to  $T$ -odd pion fragmentation

$$H_1^\perp(z, k_\perp) = \frac{N'^2 f^2 g^2}{(2\pi)^4} \frac{1}{4z} \frac{(1-z)}{z} \frac{\mu}{\Lambda'(k_\perp^2)} \frac{M_\pi}{k_\perp^2} \mathcal{R}(z, k_\perp^2)$$

$$\text{where, } \Lambda'(k_\perp^2) = k_\perp^2 + \frac{1-z}{z^2} M_\pi^2 + \frac{\mu^2}{z} - \frac{1-z}{z} m^2$$

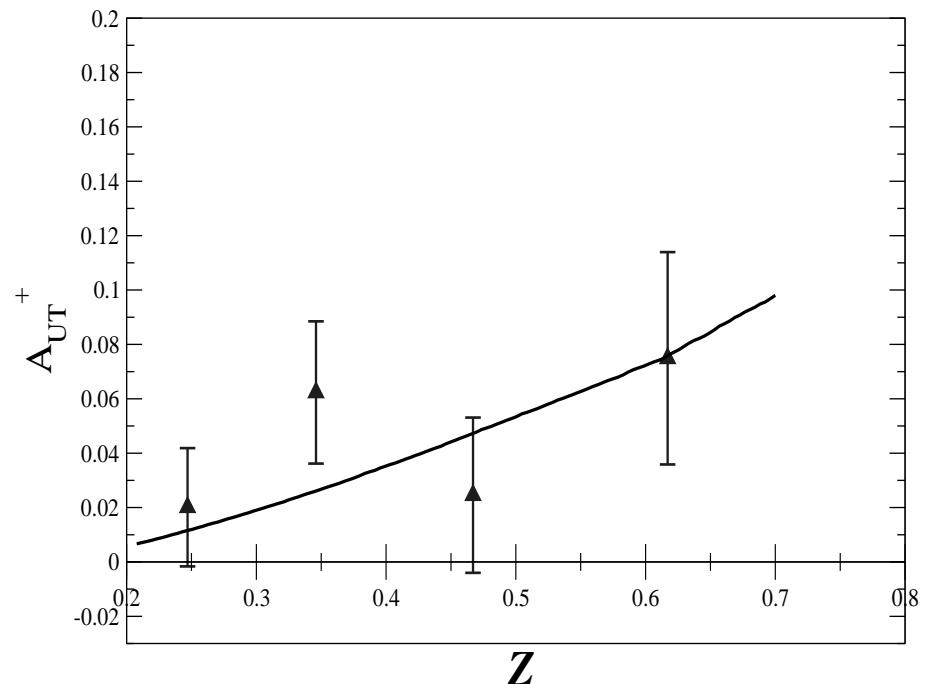
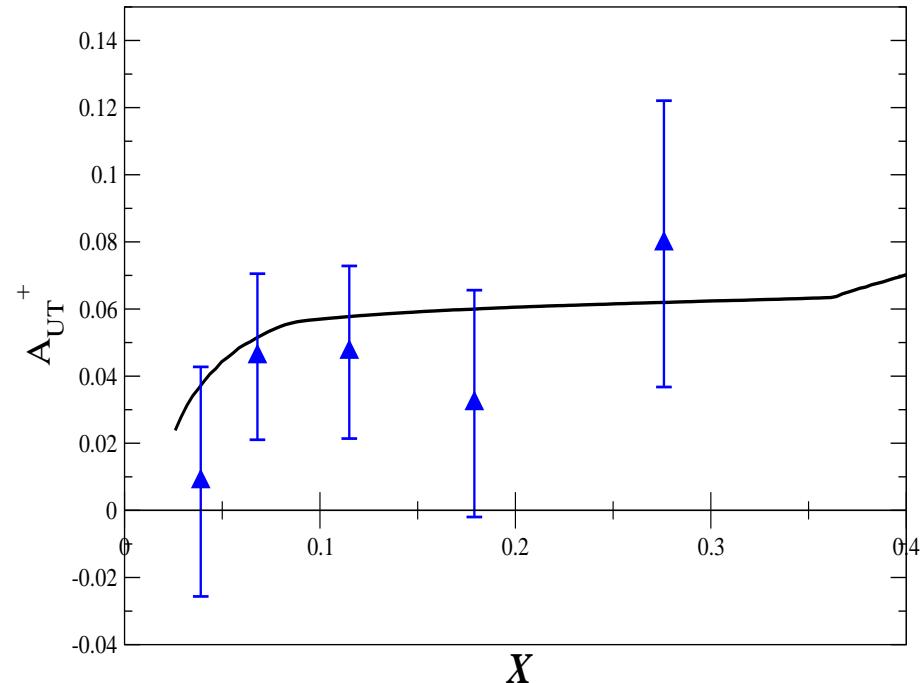
# Collins Asymmetry

Gamberg, Goldstein, Oganessyan PRD 2003: updated For the HERMES kinematics

$1 \text{ GeV}^2 \leq Q^2 \leq 15 \text{ GeV}^2$ ,  $4.5 \text{ GeV} \leq E_\pi \leq 13.5 \text{ GeV}$ ,  $0.2 \leq x \leq 0.41$ ,  $0.2 \leq z \leq 0.7$ ,  $0.2 \leq y \leq 0.8$ ,  $\langle P_{h\perp}^2 \rangle = 0.25 \text{ GeV}^2$

$$\langle \frac{P_{h\perp}}{M_\pi} \sin(\phi + \phi_s) \rangle_{UT} = |S_T| \frac{2(1-y) \sum_q e_q^2 h_1(x) z H_1^{\perp(1)}(z)}{(1+(1-y)^2) \sum_q e_q^2 f_1(x) D_1(z)}.$$

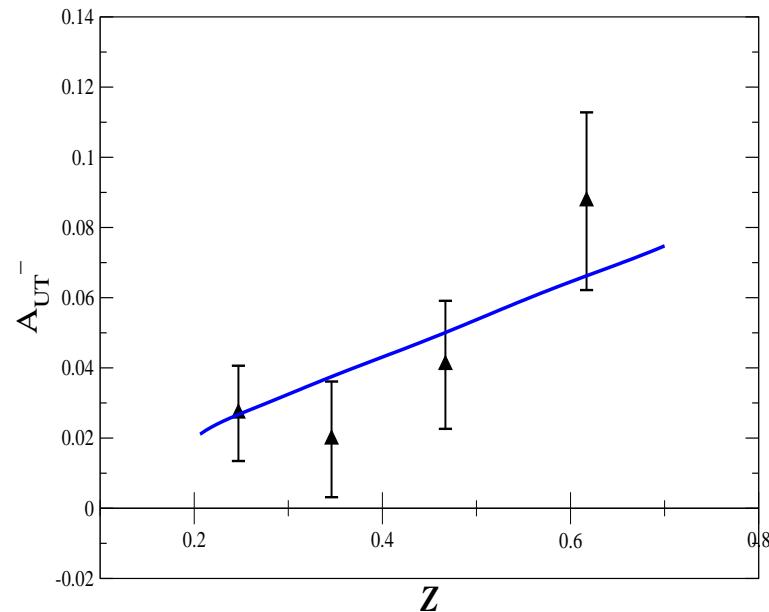
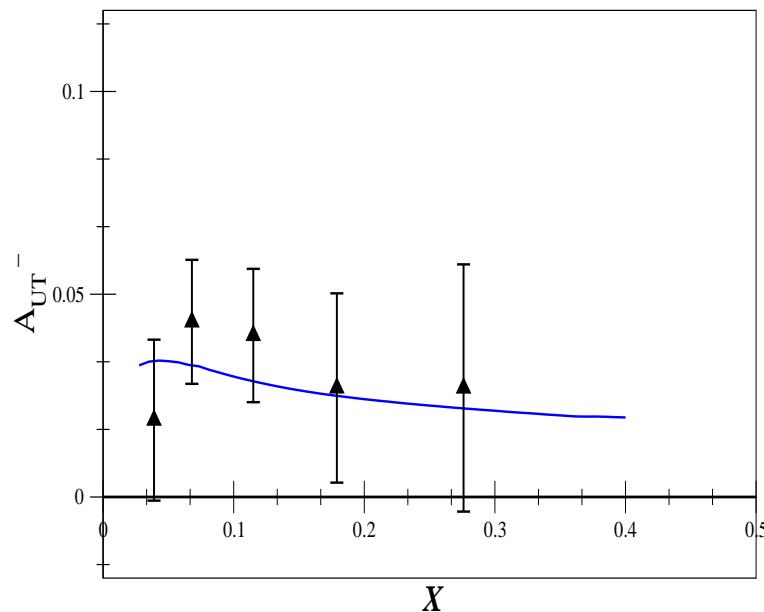
Data from A. Airapetian et al. PRD 94, 2005



# Estimates for Sivers Asymmetry

Data from A. Airapetian et al. PRL94,2005

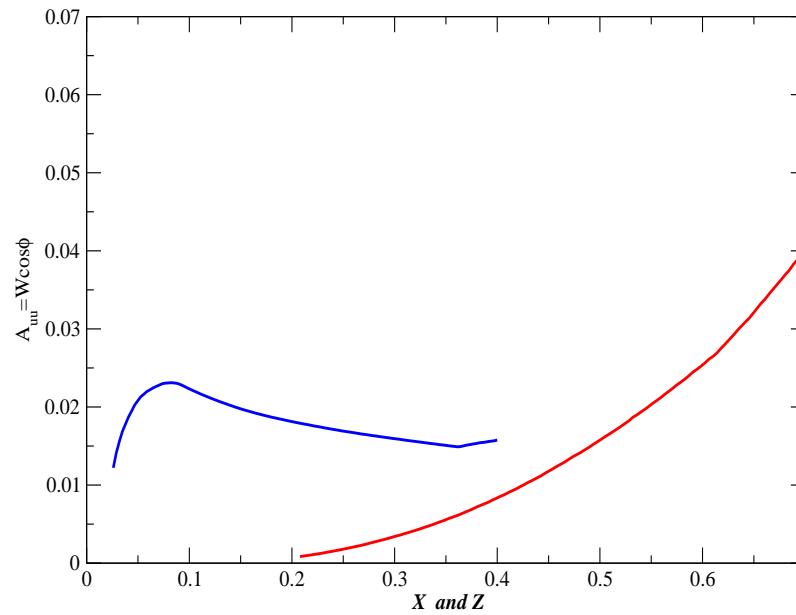
$$\begin{aligned}
 \left\langle \frac{|P_{h\perp}|}{M} \sin(\phi - \phi_S) \right\rangle_{UT} &= \frac{\int d^2 P_{h\perp} \frac{|P_{h\perp}|}{M} \sin(\phi - \phi_S) d\sigma}{\int d^2 P_{h\perp} d\sigma} \\
 &= \frac{(1 + (1 - y)^2) \sum_q e_q^2 f_1^{\perp(1)}(x) z D_1^q(z)}{(1 + (1 - y)^2) \sum_q e_q^2 f_1(x) D_1(z)},
 \end{aligned}$$



## Double T-odd $\cos 2\phi$ asymmetry

Transversity of quarks inside an unpolarized hadron, and  $\cos 2\phi$  asymmetries in unpolarized semi-inclusive DIS

$$\langle \frac{|P_{h\perp}^2|}{MM_h} \cos 2\phi \rangle_{UU} = \frac{\int d^2 P_{h\perp} \frac{|P_{h\perp}^2|}{MM_h} \cos 2\phi d\sigma}{\int d^2 P_{h\perp} d\sigma} = \frac{8(1-y) \sum_q e_q^2 h_1^{\perp(1)}(x) z^2 H_1^{\perp(1)}(z)}{(1+(1-y)^2) \sum_q e_q^2 f_1(x) D_1(z)}$$

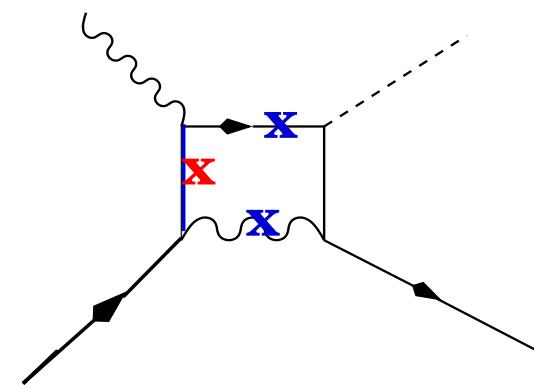
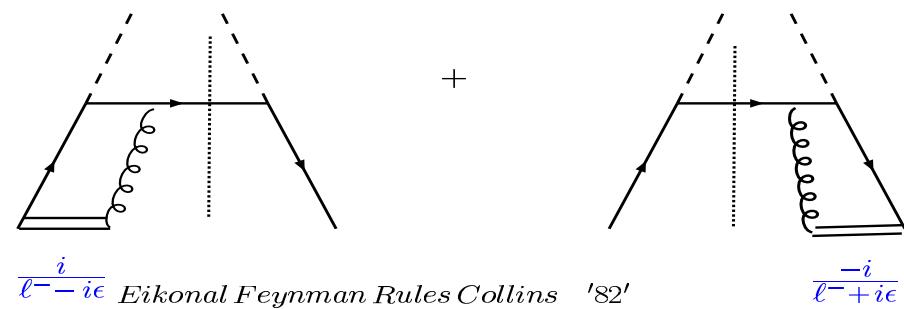
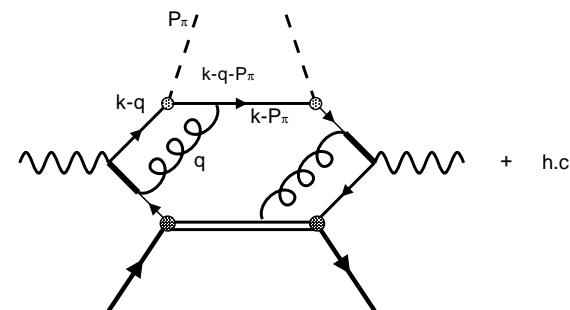
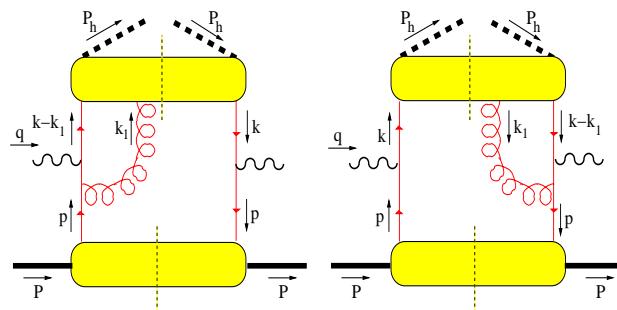


## Spectator Model: Gauge Link Contribution to Collins Function

Metz: PBL 2002, Gamberg, Goldstein, Oganessyan PRD: 2003: Bacchetta, Metz, Jang: PLB: 2003, Amrath, Bacchetta, Metz:

PRD 2005, G.G. in progress

$$\Delta^{[\sigma^\perp - \gamma_5]}(z, k_\perp) = \frac{1}{4z} \int dk^+ \text{Tr}(\gamma^- \gamma^\perp \gamma_5 \Delta) \Big|_{k^- = P_\pi^- / z} \quad \text{Boer, Pijlman, Muders: NPB 2003}$$



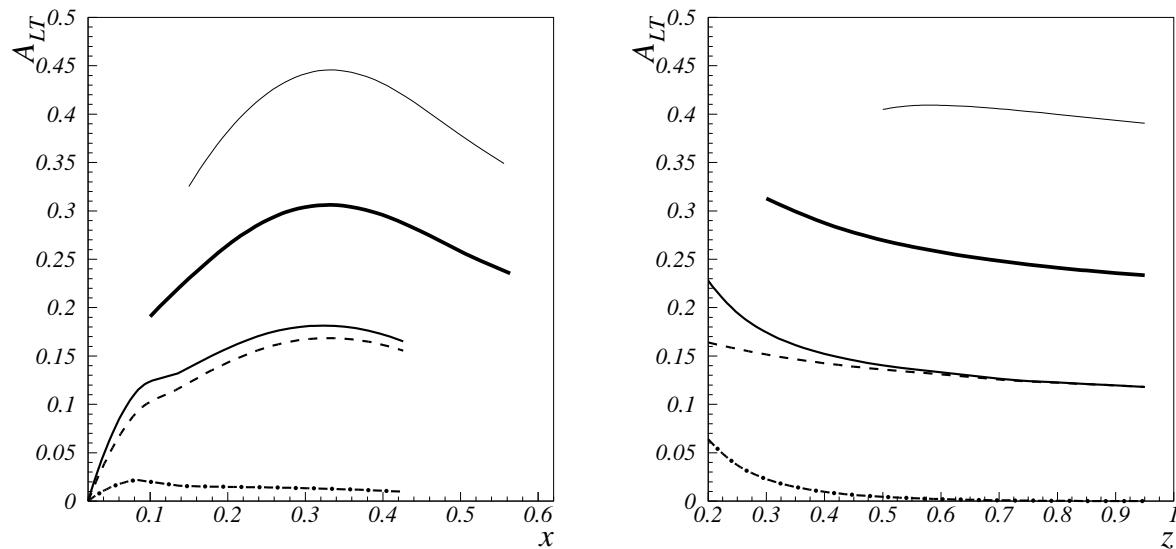
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PRD 2005, Gamberg Goldstein in progress

- Is the eikonal pole in the physical regime of the Collins function Correlator? and or off shell  
 $\gamma + q \rightarrow \pi + q'$
- Explore Pole Structure of Loop Integral
  - ★ Using Cauchy's theorem to evaluate the Color Gauge invariant Correlator eikonal pole exists at  $z = 1$ : exclusive limit.
  - ★ Can we deform the pole away if take into account  $\ell_\perp$ ?
  - ★ Evaluate the box diagram taking the eikonal limit on the fragmenting quark but keeping a mass correction: eikonal pole outside the physical regime ie  $\ell^- < 0$
  - ★ Evaluation of Cuts in *S – channel* w/o mass correction indicates L.C. divergence  $\rightarrow \ell^- \rightarrow 0$ : if regulate, cancels.
  - ★ ? Consistent with Correlator definition? Yes “maybe”....
  - ★ In correlator keep  $n$  off light cone  $n \cdot A$ ,  $n = (n^-, n^+)$  (see Ji, Yuan, Ma PLB: 2004)
  - ★ Pick up poles contributions on fragmenting quark and gluon  $\Rightarrow$  equivalent to cut in *S*-channel of box.
  - ★ Within spectator model, em suggests Collins Function universal between  $e^+e^-$  and SIDIS.
- Earlier model calculation based color Gauge Inv. Correlator definition and transcription of fragmentation functions from distribution functions in spectator model: **Jakob, Mulders, Rodrigues NPB: 1997, Bacchetta, Boffi, Jakob EJPC 2000**.

- \* SIDIS: Jaffe and Ji PRL:1993 encountered at twist three level Estimate of this effect, Gamberg, Hwang, Oganessyan PLB:2004

$$A_{LT} = \frac{\lambda_e |S_T| \sqrt{1-y} \frac{4}{Q} \left[ M x g_T(x) D_1(z) + M_h h_1(x) \frac{E(z)}{z} \right]}{\frac{[1+(1-y)^2]}{y} f_1(x) D_1(z)}$$

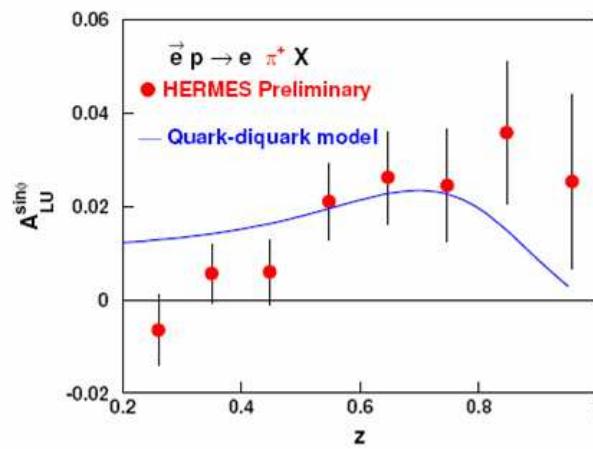
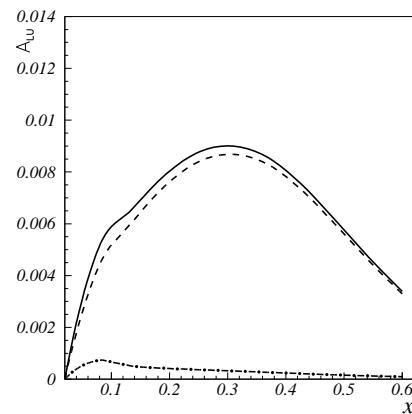


$A_{LT}$  for  $\pi^+$  production function of  $x$  and  $z$  at 27.5 GeV energy. The dashed and dot-dashed curves correspond contributions of two terms above respectively, and full curve is sum. Thin curve corresponds to 6 GeV and the thick to 12 GeV energies.

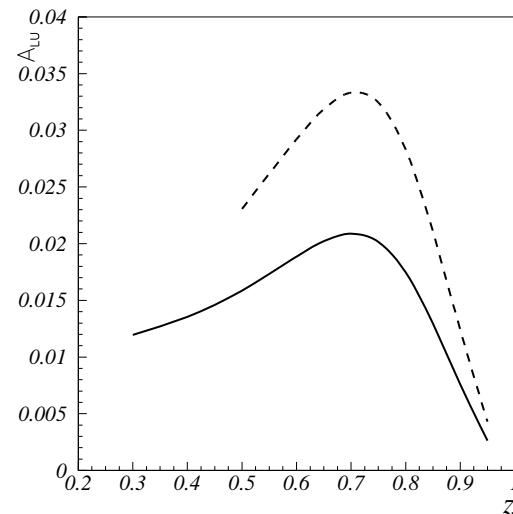
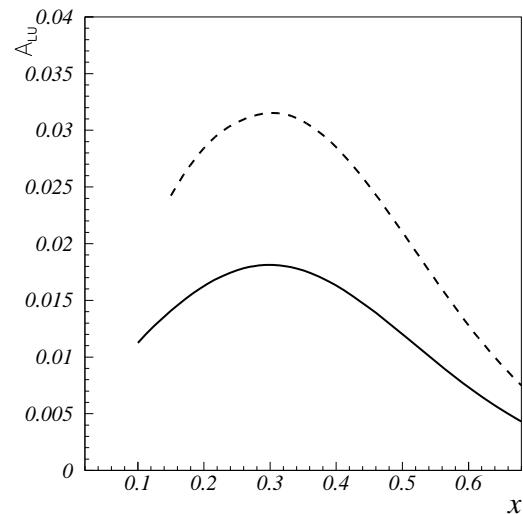
★ Bean Asymmetry Estimate of this effect, Gamberg, Hwang, Oganessyan PLB:2004

F. Yuan, PLB: 2004. Metz and Schleigel, hep-ph/0403182, Bacchetta *et al* hep-ph/0405154. extra terms  $g^\perp$  ? how big?

$$\langle |P_{h\perp}| \sin \phi \rangle_{LU} = \lambda_e \sqrt{1-y} \frac{4}{Q} M M_h \left[ x e(x) z H_1^{\perp(1)}(z) + h_1^{\perp(1)}(x) E(z) \right],$$



$A_{LU}$  for  $\pi^+$  production as a function of  $x$  and  $z$  at 27.5 GeV energy. Dashed and dot-dashed curves correspond first and second terms above respectively, and full curve, the sum.



JLAB Kin. 6 and 12 GeV

# SUMMARY

- Going beyond the collinear approximation in PQCD recent progress has been achieved in characterizing transverse SSA and azimuthal asymmetries in terms of absorptive scattering.
- Central to this understanding is the role that transversity properties of quarks and hadrons assume in terms of correlations between transverse momentum and transverse spin in QCD hard scattering.
- These asymmetries provide a window to explore novel quark distribution and fragmentation functions which constitute essential information about the spin, transversity and generalized momentum structure of hadrons.
- Along with the chiral odd transversity  $T$ -even distribution function, existence of  $T$ -odd distribution and fragmentation functions can provide an explanation for the substantial asymmetries that have been observed in inclusive and semi-inclusive scattering reactions.
- We consider the angular correlations in SDIS and Drell Yan from the standpoint of “rescattering” mechanism which generate  $T$ -odd, intrinsic transverse momentum,  $k_\perp$ , dependent *distribution and fragmentation* functions at leading twist
- We have evaluated  $T$ -odd contributions to azimuthal and SSA and modeled intrinsic  $k_\perp$  with Gaussian “regularization” in  $\langle k_\perp \rangle$
- ★ Azimuthal asymmetries in Drell Yan and SSA measured at HERMES and COMPASS, JLAB, Belle, GSI-PAX *may reveal* the extent to which these leading twist T-odd effects are generating the data