

# SSA in SIDIS & DY

Peter Schweitzer

Institute for Theoretical Physics II, Ruhr-University Bochum, Germany

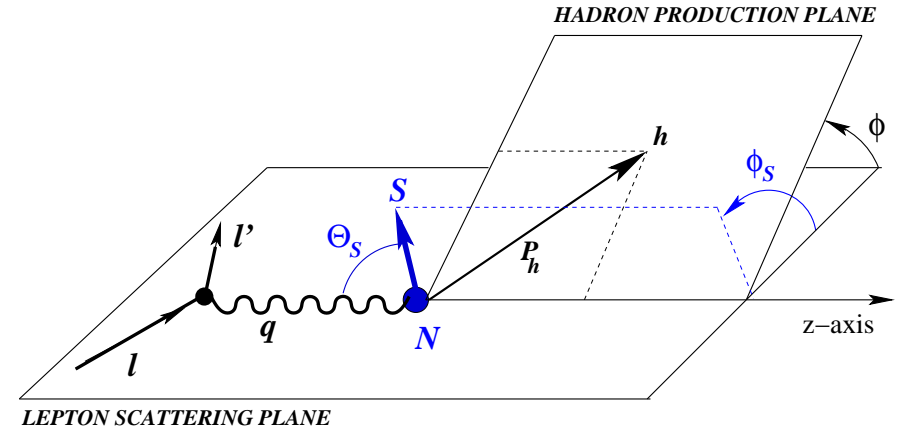
in collaboration with *J. C. Collins, A. V. Efremov, K. Goeke, M. Grosse Perdekamp, S. Menzel, A. Metz*

## Overview:

- Sivers effect in SIDIS
- *Preliminary* HERMES data on  $A_{UT}^{\sin(\phi-\phi_S)P_{h\perp}/M_N}$  & *final* HERMES data on  $A_{UT}^{\sin(\phi-\phi_S)}$ 
  - Use & test(!) large  $N_c$  (= number of colours) QCD limit for Sivers function
  - Gauss model for distribution of transverse parton momenta
  - COMPASS (deuteron) vs. HERMES (proton) data
- Drell-Yan process
- Conclusions

## HERMES & COMPASS data

$$\frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow} \propto \underbrace{\dots \sin(\phi - \phi_S)}_{\text{Sivers}} + \underbrace{\dots \sin(\phi + \phi_S)}_{\text{Collins effect}}$$



- **Sivers effect:** first *unambiguous* evidence for “naively/artificially T-odd” distribution in SIDIS
- Sivers SSA  $\propto \sum_a e_a^2 f_{1T}^{\perp a} D_1^a$  and  $D_1^a(z)$  known, e.g. Kretzer, Leader, Christova, EPJC 22 (2001) 269
- $f_{1T}^{\perp a}$  only unknown, i.e. possible to extract direct information on Sivers function
- subject to considerable interest in literature:
  - Efremov, Goeke, Menzel, Metz and PS, PLB 612 (2005) 233, hep-ph/0412353
  - Anselmino, Boglione, D’Alesio, Kotzinian, Murgia and Prokudin, PRD 71 (2005) 074006, hep-ph/0501196
  - Anselmino, Boglione, D’Alesio, Kotzinian, Murgia and Prokudin, hep-ph/0507181
  - Vogelsang and Yuan, hep-ph/0507266
  - Collins, Efremov, Goeke, Menzel, Metz and PS, hep-ph/0509076
- **Collins effect**  $\propto \sum_a e_a^2 h_1^a H_1^{\perp a}$  equally interesting, first *unambiguous* evidence in SIDIS. Data indicate fascinating properties  $\rightarrow$  favoured vs. unfavoured  $H_1^\perp$ .  
**But:** Two unknowns, more difficult. See Vogelsang and Yuan, *op. cit.*; Efremov et al., talk at SIR’05

- FINAL HERMES & COMPASS data [PRL 94 \(2005\) 012002](#) & [PRL 94 \(2005\) 202002](#)

$$A_{UT}^{\sin(\phi-\phi_S)} \equiv \frac{\int d\phi \sin(\phi - \phi_S)(N^\uparrow - N^\downarrow)}{\frac{1}{2} \int d\phi (N^\uparrow + N^\downarrow)}$$

- “preferable” PRELIMINARY HERMES data [Makins, “Transversity Workshop”, Oct. 2003, Athens.](#)  
[Seidl, DIS’2004, April 2004, Štrbské Pleso.](#) [Gregor, Acta Phys. Pol. B 36 \(2005\) 209](#)

$$A_{UT}^{\sin(\phi-\phi_S)P_{h\perp}/M_N} \equiv \frac{\int d\phi \sin(\phi - \phi_S)P_{h\perp}/M_N(N^\uparrow - N^\downarrow)}{\frac{1}{2} \int d\phi (N^\uparrow + N^\downarrow)} = (-2) \frac{\sum_a e_a^2 x f_{1T}^{\perp(1)a}(x) z D_1^{a/\pi}(z)}{\sum_a e_a^2 x f_1^a(x) D_1^{a/\pi}(z)}$$

with  $f_{1T}^{\perp(1)a}(x) \equiv \int d^2\mathbf{p}_T \frac{\mathbf{p}_T^2}{2M_N^2} f_{1T}^{\perp a}(x, \mathbf{p}_T^2)$  [Boer and Mulders, PRD 57 \(1998\) 5780](#)

- In which sense is  $A_{UT}^{\sin(\phi-\phi_S)}$  “not preferable”? See later in detail.
- **WARNING!** Keep in mind: In any case **approximation** – neglect soft factors (!)  
[Ji, Ma and Yuan, PRD 71 \(2005\) 034005](#), [PLB 597 \(2004\) 299](#). [Collins and Metz, PRL 93 \(2004\) 252001](#).

$A_{UT}^{\sin(\phi-\phi_S)P_{h\perp}/M_N}$   $\xleftrightarrow{\text{directly}}$   $f_{1T}^{\perp(1)a}(x)$  explored by Efremov *et al.*, PLB 612 (2005) 233, hep-ph/0412353

- error bars sizeable  $\rightarrow$  minimize # of free parameters in fit Ansatz  $\rightarrow$  **theoretical constraints**

I use large- $N_c$  limit  $\boxed{f_{1T}^{\perp u} = -f_{1T}^{\perp d}}$  modulo  $1/N_c$  corrections Poylyitsa hep-ph/0301236 (★)

II satisfy sum rule  $\sum_{a=g,u,d,\dots} \int dx f_{1T}^{\perp(1)a}(x) = 0$  Burkardt, PRD 69 (2004) 057501 and 091501

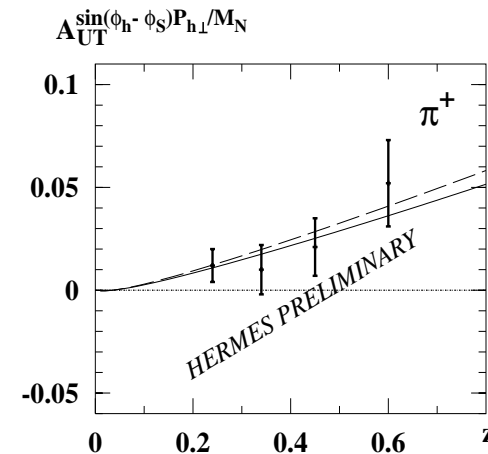
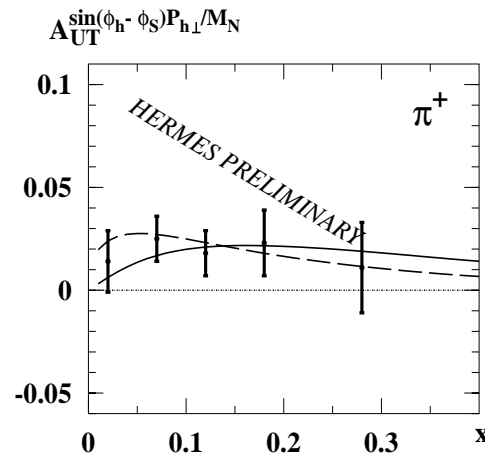
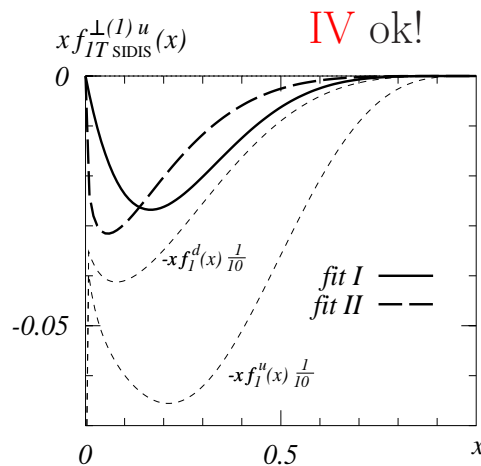
III explore “connection” of  $f_{1T}^{\perp a}$  to GPD,  $E^q(x, 0, 0) \propto (1-x)^4$  or  $(1-x)^5$  Brodsky *et al.*, Burkardt

IV respect positivity  $\frac{|\mathbf{p}_T|}{M_N} |f_{1T}^{\perp a}(x, \mathbf{p}_T^2)| \leq f_1^a(x, \mathbf{p}_T^2)$  Bacchetta, *et al.*, PRL 85 (2000) 712

• fit Ansatz:  $x f_{1T}^{\perp(1)u} = -x f_{1T}^{\perp(1)d} = Ax^b(1-x)^5$  neglect  $\bar{q}, s, \dots$  Respects I, II, III!

• fit result:  $x f_{1T}^{\perp(1)u} = \begin{cases} -0.4x(1-x)^5 & \text{for } b=1 \text{ fixed} \\ -0.1x^{0.3}(1-x)^5 & \text{both } A, b \text{ free} \end{cases}$

Cross check!



★ Footnote for historical correctness:

- M. Anselmino, V. Barone, A. Drago and F. Murgia,  
“Non-standard time reversal for particle multiplets and the spin-flavor structure of hadrons,”  
Nucl. Phys. Proc. Suppl. **105** (2002) 132, hep-ph/0111044
- M. Anselmino, V. Barone, A. Drago and F. Murgia,  
“Non-standard time reversal and transverse single-spin asymmetries”,  
hep-ph/0209073
- P. V. Pobylitsa,  
“T-odd quark distributions: QCD versus chiral models”,  
hep-ph/0212027
- P. V. Pobylitsa,  
“Transverse-momentum dependent parton distributions in large- $N_c$  QCD”,  
hep-ph/0301236.
- A. Drago,  
“Time-reversal odd distribution functions in chiral models with vector mesons”,  
Phys. Rev. D **71** (2005) 057501, hep-ph/0501282

## Comments

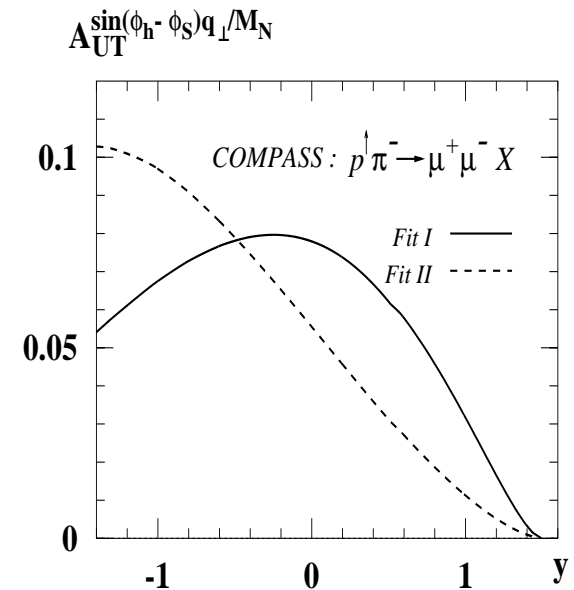
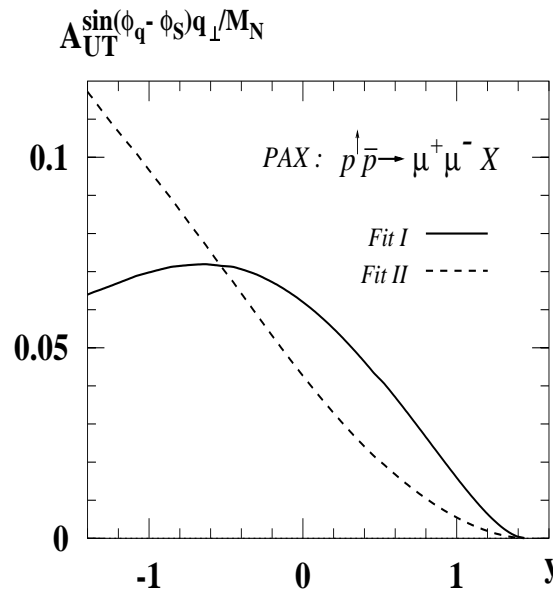
- supports picture by Burkardt, *op. cit.*  $\int dx f_{1T \text{SIDIS}}^{\perp(1)u}(x) \propto -\kappa^u < 0$ ,  $\int dx f_{1T \text{SIDIS}}^{\perp(1)d}(x) \propto -\kappa^d > 0$
- Where are  $1/N_c$  corrections? Good place to see:  $f_{1T}^{\perp u/\text{deuteron}} \approx f_{1T}^{\perp u/p} + f_{1T}^{\perp u/n} \approx f_{1T}^{\perp u} + f_{1T}^{\perp d}$
- COMPASS  $A_{UT, \text{deuteron}}^{\sin(\phi - \phi_s)} \sim 1/N_c \sim 0$  within error bars
- Thus HERMES & COMPASS compatible,  $f_{1T}^{\perp a}|_{\text{SIDIS}}$  obeys large  $N_c$  predictions

## Application

- important prediction  
Collins, PLB 536 (2002) 43

$$f_{1T}^{\perp a}|_{\text{SIDIS}} = - f_{1T}^{\perp a}|_{\text{DY}}$$

- check in experiment!
- PAX at GSI  $p^\uparrow \bar{p} \rightarrow \mu^+ \mu^- X$
- COMPASS  $p^\uparrow \pi^- \rightarrow \mu^+ \mu^- X$
- possible to confirm (or ...)



**Complication** at the present stage: Status of used HERMES data PRELIMINARY.

Not recommended to use these data, because analysis of possible error correlations not finished

HERMES Collaboration, private communication

Necessary to solidify our conclusions → study the FINAL but “less preferably” weighted data.

$$A_{UT}^{\sin(\phi_h - \phi_S)} = (-2) \frac{\sum_a e_a^2 \int d^2 \mathbf{P}_{h\perp} \int d^2 \mathbf{p}_T \int d^2 \mathbf{K}_T \sin(\phi_h - \phi_S) \sin(\phi_{\mathbf{p}_T} - \phi_S) \frac{|\mathbf{p}_T|}{M_N} \delta^{(2)}(\mathbf{p}_T - \mathbf{K}_T - \mathbf{P}_{h\perp}/z) x f_{1T}^{\perp a}(x, \mathbf{p}_T^2) D_1^a(z, \mathbf{K}_T^2)}{\sum_b e_b^2 x f_1^b(x) D_1^b(z)}$$

→ Have to assume some **model** for transverse parton momenta to solve convolution integrals.

Already committed a crime: soft factors →  $\delta^{(2)}(\dots)$ . Why not another?

Let us use **Gaussian model**

$$\begin{aligned} f_1^a(x, \mathbf{p}_T^2) &\equiv f_1^a(x) \frac{\exp(-\mathbf{p}_T^2/p_{\text{unp}}^2)}{\pi p_{\text{unp}}^2}, \\ f_{1T}^{\perp a}(x, \mathbf{p}_T^2) &\equiv f_{1T}^{\perp a}(x) \frac{\exp(-\mathbf{p}_T^2/p_{\text{Siv}}^2)}{\pi p_{\text{Siv}}^2} \implies A_{UT}^{\sin(\phi - \phi_S)} = (-2) \frac{a_{\text{Gauss}} \sum_a e_a^2 x f_{1T}^{\perp(1)a}(x) D_1^a(z)}{\sum_b e_b^2 x f_1^b(x) D_1^b(z)}, \\ D_1^a(z, \mathbf{K}_T^2) &\equiv D_1^a(z) \frac{\exp(-\mathbf{K}_T^2/K_{D_1}^2)}{\pi K_{D_1}^2} \quad \text{with } a_{\text{Gauss}} = \frac{\sqrt{\pi}}{2} \frac{M_N}{\sqrt{p_{\text{Siv}}^2 + K_{D_1}^2/z^2}}. \end{aligned}$$

Mulders and Tangerman, NPB 461 (1996) 197. Efremov, Goeke and PS, PLB 568 (2003) 63.

Price to pay: 2 free parameters  $p_{\text{unp}}^2$ ,  $\underbrace{p_{\text{Siv}}^2, K_{D_1}^2}_{\text{How to fix?}}$  if we assume them flavour- and  $x$ - or  $z$ -independent.

Note: In numerous processes reasonable **model**, if: transv. momenta  $\ll$  hard scale

D’Alesio and Murgia, PRD 70 (2004) 074009.

- In HERMES experiment  $\langle P_{h\perp} \rangle \sim 0.4 \text{ GeV} \ll \sqrt{\langle Q^2 \rangle} \sim 1.5 \text{ GeV}$ .

- more precisely:  $\langle P_{h\perp}(z) \rangle$  of hadrons produced from a deuterium target [HERMES, PLB 562 \(2003\) 182](#)

- $\langle P_{h\perp}(z) \rangle \stackrel{\text{Gauss}}{=} \frac{\sqrt{\pi}}{2} \sqrt{z^2 p_{\text{unp}}^2 + K_{D_1}^2}$

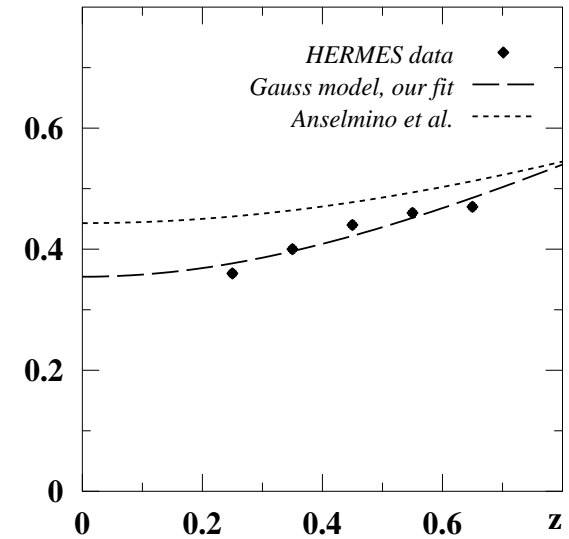
- $p_{\text{unp}}^2 = 0.33 \text{ GeV}^2$  and  $K_{D_1}^2 = 0.16 \text{ GeV}^2 \rightarrow \text{Ok!}$

- compare to study of Cahn effect by [Anselmino et al.](#) where

$$p_{\text{unp}}^2 = 0.2 \text{ GeV}^2 \text{ and } K_{D_1}^2 = 0.25 \text{ GeV}^2 \text{ obtained.}$$

$\implies$  Good qualitative agreement!

$\langle \mathbf{P}_{h\perp}(z) \rangle$  in GeV



**Ok,  $p_{\text{unp}}^2$ ,  $K_{D_1}^2$  fixed. But what is  $p_{\text{Siv}}^2$  ?**

- Of course, cannot be zero! Otherwise  $f_{1T}^\perp(x, \mathbf{p}_T^2) \rightarrow f_{1T}^\perp(x) \delta^{(2)}(\mathbf{p}_T) \implies A_{UT}^{\sin(\phi - \phi_S)} \rightarrow 0$

- Also, cannot be arbitrarily large. To satisfy positivity in Gauss Ansatz

$$p_{\text{Siv}}^2 \leq \frac{p_{\text{unp}}^2}{1 + \frac{2M_N^2 p_{\text{unp}}^2}{e p_{\text{Siv}}^4} \left( \frac{f_{1T}^{\perp(1)a}(x)}{f_1^a(x)} \right)^2}.$$

Parameter restricted  $0 < p_{\text{Siv}}^2 < p_{\text{unp}}^2$ . Admittedly vague. But, which is astonishing, ...



- $0 < p_{\text{Siv}}^2 < 0.33 \text{ GeV}^2$  **sufficient**, if we extract  $f_{1T}^{\perp(1)a}(x) = \frac{p_{\text{Siv}}^2}{2M_N^2} f_{1T}^{\perp a}(x)$  and not  $f_{1T}^{\perp a}(x)$  itself.

$$\Rightarrow \text{needed only in } a_{\text{Gauss}} = \frac{\sqrt{\pi}}{2} \frac{M_N}{\sqrt{p_{\text{Siv}}^2 + K_{D_1}^2/z^2}} \Rightarrow 0.72 < a_{\text{Gauss}} < 0.83 \Rightarrow 10\% \text{ uncertainty!}$$

## Procedure

- Ansatz  $x f_{1T}^{\perp(1)u}(x) = -x f_{1T}^{\perp(1)d}(x) = Ax^b(1-x)^5$ , neglect anti- and heavier quarks.
- Choose some  $p_{\text{Siv}}^2 \in [0, 0.33 \text{ GeV}^2]$
- Fit to *FINAL* HERMES data

- Check (within  $2\text{-}\sigma$ ) positivity  $p_{\text{Siv}}^2 \leq \frac{p_{\text{unp}}^2}{1 + \frac{2M_N^2 p_{\text{unp}}^2}{e p_{\text{Siv}}^4} \left( \frac{f_{1T}^{\perp(1)a}(x)}{f_1^a(x)} \right)^2}$

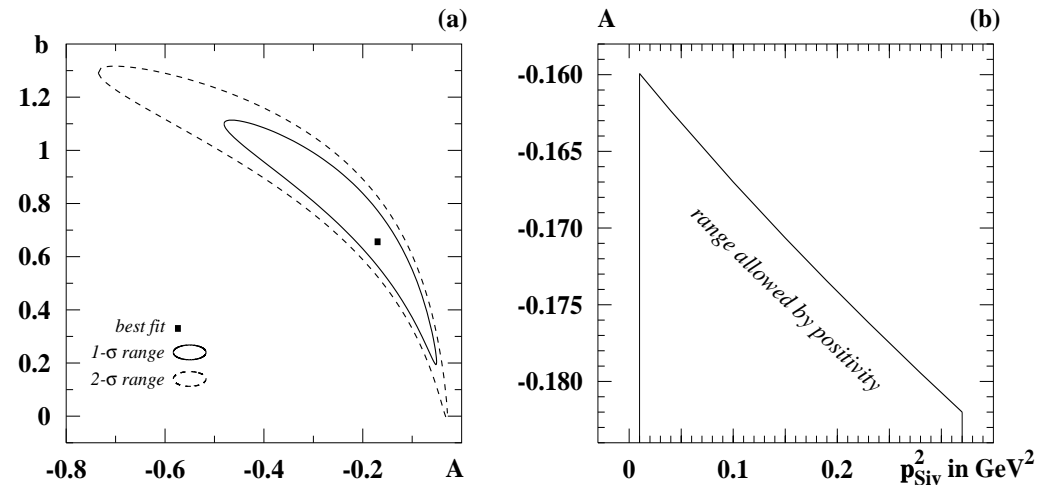
## Result

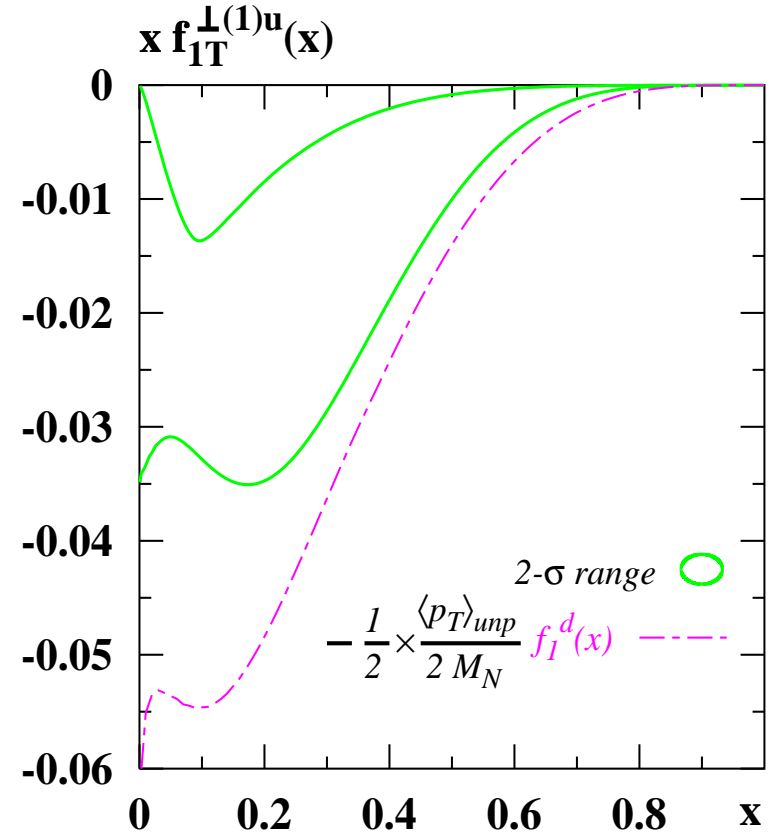
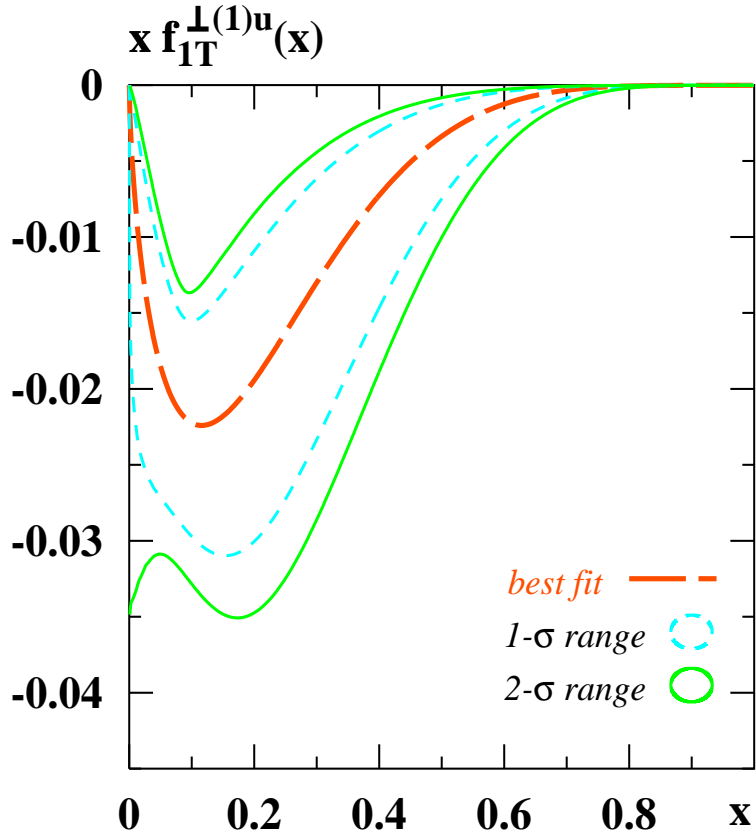
- Best fit ( $\chi^2 \sim 2.2$ )

$$x f_{1T}^{\perp(1)u}(x) = -(0.16 \dots 0.18) x^{0.66} (1-x)^5$$

$$\text{for } p_{\text{Siv}}^2 = (0.01 \dots 0.32) \text{ GeV}^2$$

- refers to scale  $\sim 2.5 \text{ GeV}^2$



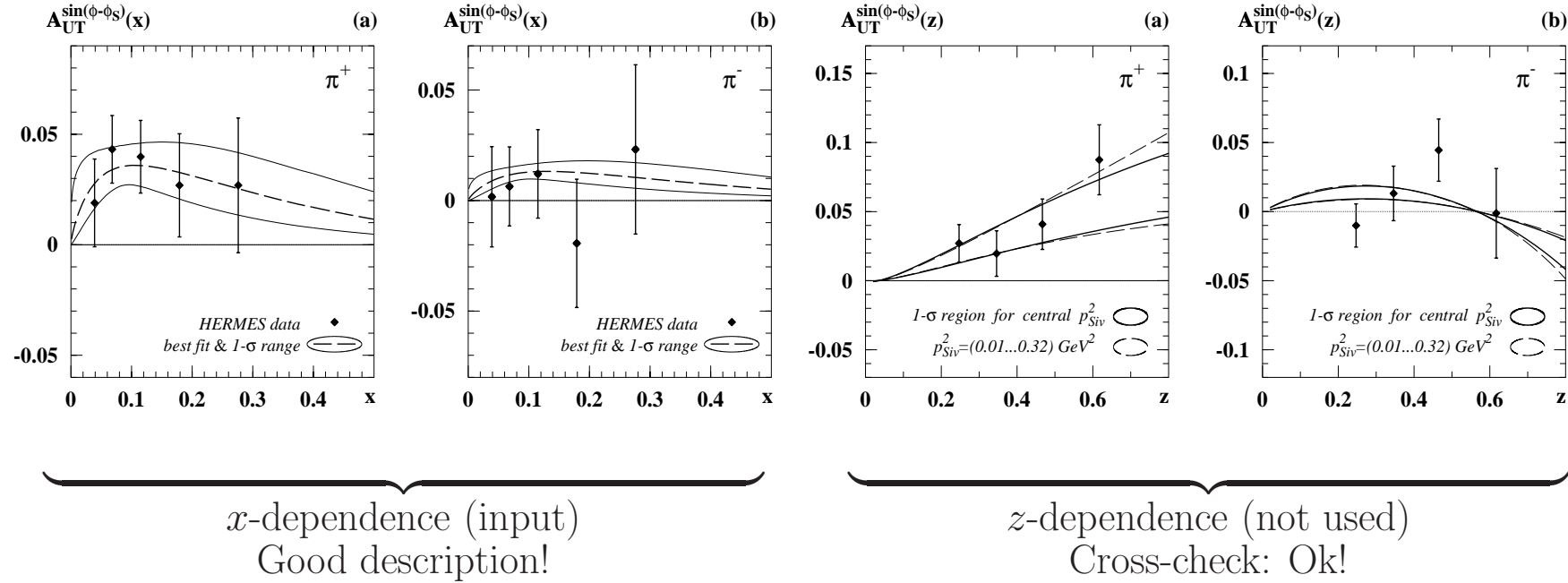


Positivity ok! & Not small!!!

- In large- $N_c$ :  $\frac{|\mathbf{p}_T|}{M_N} |f_{1T}^{\perp u}(x, \mathbf{p}_T^2)| = \frac{|\mathbf{p}_T|}{M_N} |f_{1T}^{\perp d}(x, \mathbf{p}_T^2)| \leq f_1^u(x, \mathbf{p}_T^2) = f_1^d(x, \mathbf{p}_T^2)$
- In nature  $f_1^u(x, \mathbf{p}_T^2) > f_1^d(x, \mathbf{p}_T^2) \Rightarrow$  for  $d$ -quark stronger. To check multiply by  $|\mathbf{p}_T|$  & integrate

$$|f_{1T}^{\perp(1)u,d}(x)| \leq \frac{\langle p_T \rangle_{\text{unp}}}{2M_N} f_1^d(x) \quad \langle p_T \rangle_{\text{unp}} = \frac{\int d^2 \mathbf{p}_T |\mathbf{p}_T| f_1^q(x, \mathbf{p}_T^2)}{\int d^2 \mathbf{p}_T f_1^q(x, \mathbf{p}_T^2)} \stackrel{\text{Gauss}}{=} \frac{\sqrt{\pi}}{2} p_{\text{unp}} = 0.51 \text{ GeV}$$

How does it describe the HERMES data? Well!



Good description. However, two questions arise:

- $f_1^{\bar{d}}(x)/f_1^d(x) \sim 25\%$  at small HERMES- $x$ . Justified to neglect Sivers- $\bar{q}$  ?
- Where are  $1/N_c \sim 30\%$  corrections?

## Sivers- $\bar{q}$

- Let us simulate

$$f_{1T}^{\perp(1)\bar{q}}(x) = \underbrace{f_{1T}^{\perp(1)q}(x)}_{\text{best fit}} \times \begin{cases} \pm 25\% \\ \pm \frac{(f_1^{\bar{u}} + f_1^{\bar{d}})(x)}{(f_1^u + f_1^d)(x)} \end{cases}$$

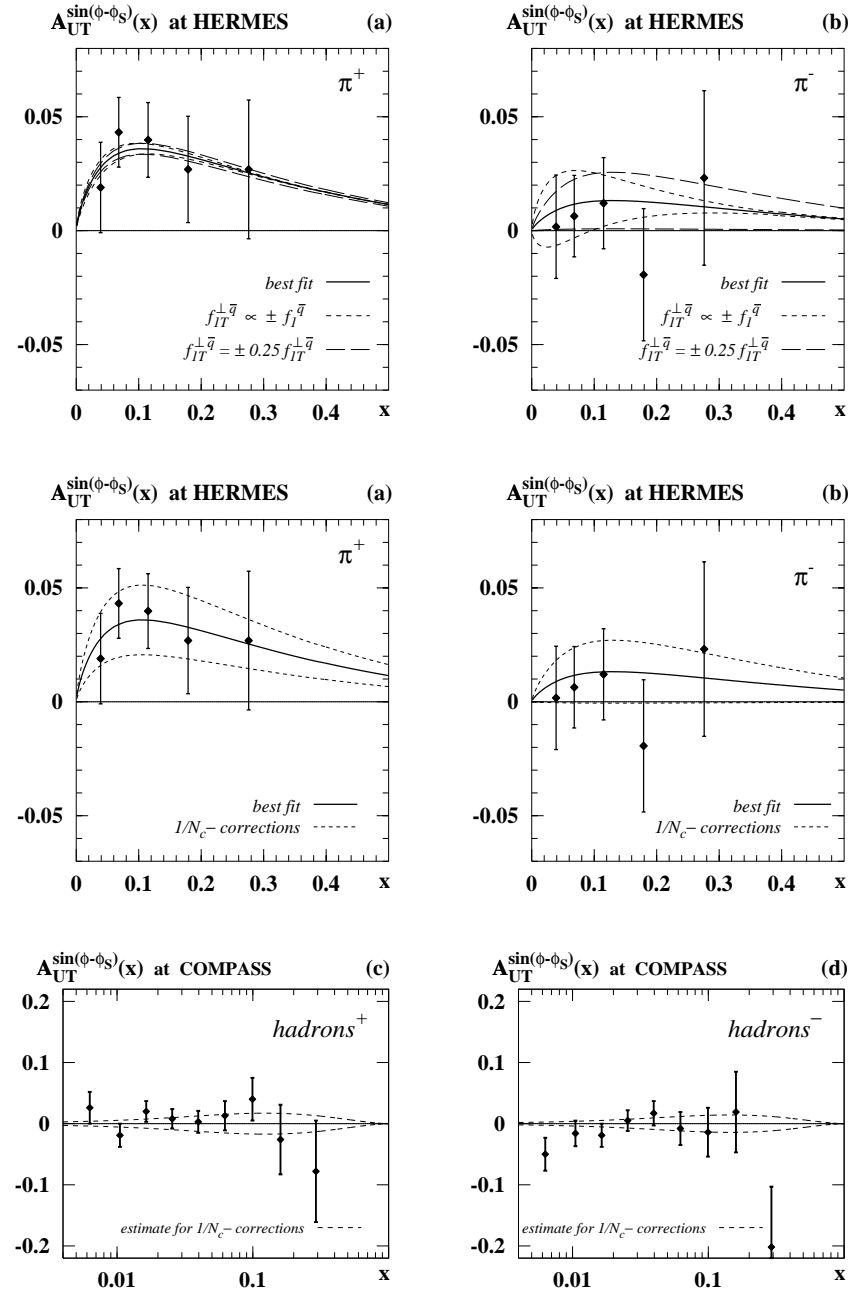
## $1/N_c$ -corrections

- $\underbrace{|(f_{1T}^{\perp u} + f_{1T}^{\perp d})(x)|}_{\mathcal{O}(N_c^2)} \ll \underbrace{|(f_{1T}^{\perp u} - f_{1T}^{\perp d})(x)|}_{\mathcal{O}(N_c^3)}$
- $(f_{1T}^{\perp u} + f_{1T}^{\perp d})(x) = \pm \frac{1}{N_c} \underbrace{(f_{1T}^{\perp u} - f_{1T}^{\perp d})(x)}_{\text{best fit}} \neq 0$

- $f_{1T}^{\perp u/\text{deuteron}} = (f_{1T}^{\perp u} + f_{1T}^{\perp d})$ , etc.

Sivers SSA on deuterium  $\stackrel{!}{=}$  "1/ $N_c$ -correction"

(Beware: Small  $x$ !)



**Conclusion:** Neglect of  $\bar{q}$ -effects and  $1/N_c$ -corrections in fit Ansatz not unjustified for present data

Maybe, **suspicion(!)**, large- $N_c$  works even *particularly well* for Siverson function because it *happens to work* particularly well for the anomalous magnetic moments “related” by  $\int dx f_{1T \text{SIDIS}}^{\perp(1)q}(x) \propto -\kappa^q$ , Burkardt, *op. cit.*

Recall:  $\kappa^u = 1.673$  and  $\kappa^d = -2.033$

In nature:  $\underbrace{|\kappa^u - \kappa^d|}_{\mathcal{O}(N_c^2)} \sim 3.706 \gg \underbrace{|\kappa^u + \kappa^d|}_{\mathcal{O}(N_c)} \sim 0.360$

**Is this an additional reason why large- $N_c$  useful for Siverson effect?** The future will show.

There are large- $N_c$  predictions also for all the other novel twist-2 functions. E.g. Boer-Mulders function

$h_1^{\perp u} = h_1^{\perp d}$  modulo  $1/N_c$  corrections, ... [Pobylitsa hep-ph/0301236](#) Worthwhile to be kept in mind.

**Where else is large- $N_c$  useful?**

Worthwhile recalling  $\frac{\Delta G(x)}{G(x)} \sim \frac{1}{N_c}$  [Efremov, Goeke and Pobylitsa, PLB 488 \(2000\) 182, hep-ph/0004196](#)

$(\Delta \bar{u} - \Delta \bar{d})(x) > |(\bar{u} - \bar{d})(x)| \neq 0$  [Diakonov, NPB 480 \(1996\) 341, hep-ph/9606314](#) Being tested ...

$f_{1T}^{\perp(1)q}$  from FINAL  
 HERMES data analysed  
 without  $P_{h\perp}$ -weight  
 hep-ph/0509076

$f_{1T}^{\perp(1)q}$  from PRELIMINARY  
 but “more preferably”  
 $P_{h\perp}$ -weighted data  
 PLB 612 (2005) 233, hep-ph/0412353

Here FINAL data on  $A_{UT}^{\sin(\phi-\phi_S)}$  **Gauss**  $\longrightarrow$  ... constrain parameters ... **etc.**  $\longrightarrow$   $f_{1T}^{\perp(1)q}(x)|_{\text{model-dependent}}$

There PRELIMINARY data on  $A_{UT}^{\sin(\phi-\phi_S)P_{h\perp}/M_N}$   $\xrightarrow{\text{directly}}$   $f_{1T}^{\perp(1)q}(x)|_{\text{“model-independent”}}$



### Does the Gauss Ansatz work?

- The same events with/without  $P_{h\perp}$ -weight!

If Gauss ok  $\implies$  the same  $f_{1T}^{\perp(1)q}(x)$

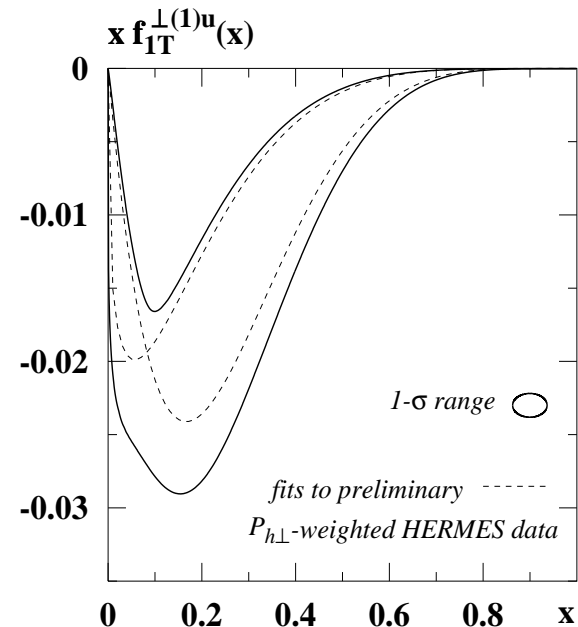
$\implies$  **Yes, it works!**

### Grain of salt:

- Be careful: Preliminary data.

“Indication”: Effects of error correlations

less dominant than statistical error of data



## Intermediate conclusion:

- Gauss Ansatz & large- $N_c$  presently useful tools
- Extraction of  $f_{1T}^{\perp(1)q}(x)$  from preliminary data “confirmed”
- Previous conclusions solidified (DY, change of sign, PAX & COMPASS, etc.)

see also [Anselmino \*et al.\*](#), and [Vogelsang and Yuan, \*op. cit.\*](#)

## How to test further Gauss model?

- Range of reliability? Limitations?
- e.g. for  $\pi^0$  (neglecting  $s$ , ...)

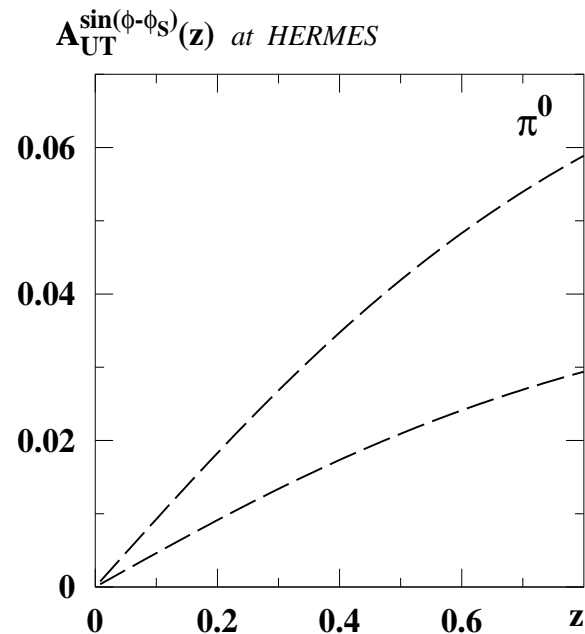
$$A_{UT}^{\sin(\phi-\phi_S)}(z) = a_{\text{Gauss}}(z) \frac{\sum_a e_a^2 \langle x f_{1T}^{\perp(1)a} \rangle}{\sum_b e_b^2 \langle x f_1^b \rangle}$$
$$\propto \frac{z}{\sqrt{\underbrace{z^2 p_{\text{Siv}}^2}_{\text{smaller}} + \underbrace{K_{D_1}^2}_{\text{larger}}}}$$

## How to see $1/N_c$ corrections?

- More precise deuteron & proton data  
most recent *PRELIMINARY* HERMES data [Diefenthaler, DIS'05, hep-ex/0507013](#) still compatible

## How to see Sivers- $\bar{q}$ ?

- Drell-Yan



# Sivers effect in DY

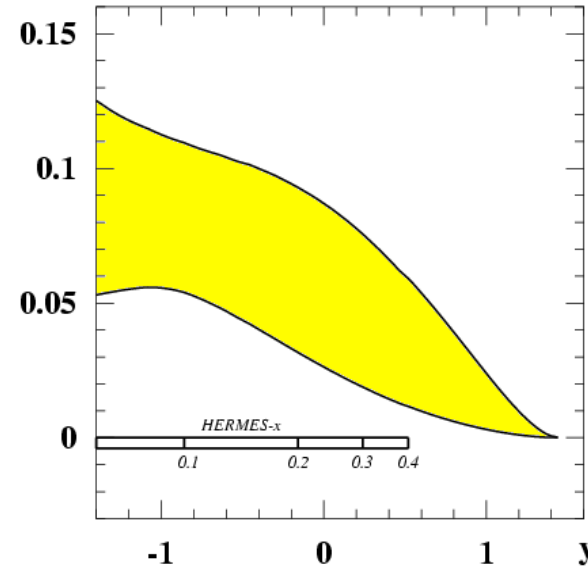
## PAX & COMPASS

- $$\left. \begin{array}{l} p^\uparrow \bar{p} \\ p^\uparrow \pi^- \end{array} \right\} \rightarrow l^+ l^- X$$

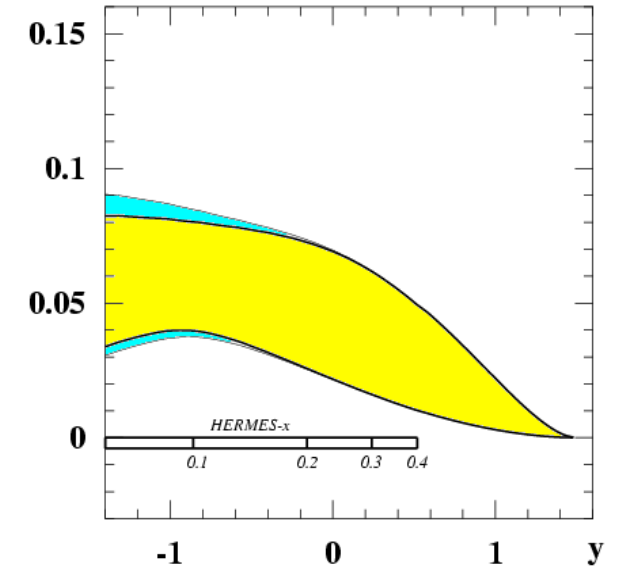
dominated by annihilations of  
valence  $q$  and valence  $\bar{q}$

$\Rightarrow$  not sensitive to Sivers- $\bar{q}$ , good!

$$A_{\text{UT}}^{\sin(\phi - \phi_S)} \text{ in } p^\uparrow \bar{p} \rightarrow l^+ l^- X \text{ at PAX}$$



$$A_{\text{UT}}^{\sin(\phi - \phi_S)} \text{ in } p^\uparrow \pi^- \rightarrow l^+ l^- X \text{ at COMPASS}$$



## RHIC

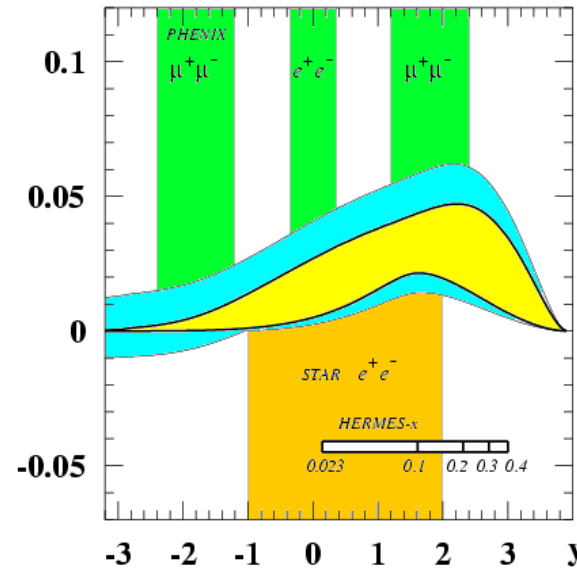
- $$p^\uparrow p \rightarrow l^+ l^- X$$

valence  $q$  and sea  $\bar{q}$  on equal footing  
sensitive to Sivers- $\bar{q}$  in certain  $y$ -region

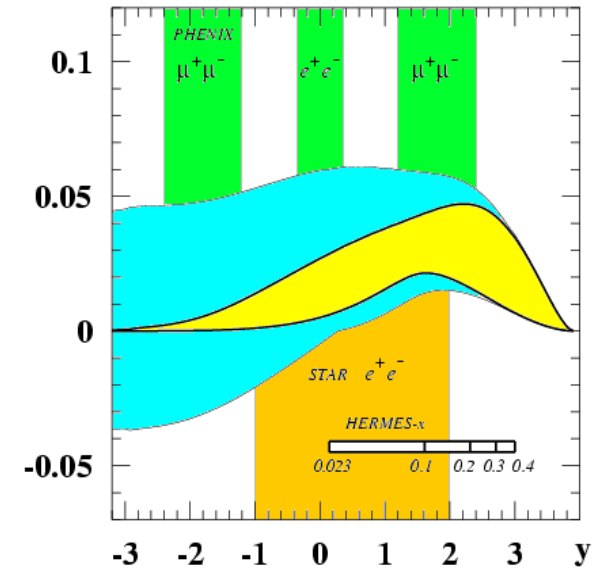
- RHIC can test “change of sign”

& provide information on Sivers- $\bar{q}$ !

$$A_{\text{UT}}^{\sin(\phi - \phi_S)} \text{ in } p^\uparrow p \rightarrow l^+ l^- X \text{ at RHIC } Q=4\text{GeV}$$



$$A_{\text{UT}}^{\sin(\phi - \phi_S)} \text{ in } p^\uparrow p \rightarrow l^+ l^- X \text{ at RHIC } Q=4\text{GeV}$$



yellow = 1- $\sigma$  region

blue = effects due to Sivers- $\bar{q}$



# Conclusions

- Previous study of preliminary HERMES data on  $A_{UT}^{\sin(\phi-\phi_S)P_{h\perp}/M_N}$  solidified by study of final HERMES data on  $A_{UT}^{\sin(\phi-\phi_S)}$
- Gaussian model  $\Rightarrow$  parameters constrained in terms of HERMES data  $\Rightarrow$  consistent with data
- Large- $N_c$  predictions used (and: tested!). Useful guideline!  
 $\Rightarrow$  naturally explains compatibility of HERMES (proton) and COMPASS (deuteron) data
- Obtained fit satisfies all presently known theoretical constraints (large- $N_c$ , Burkardt sum-rule, “connection” to GPD, positivity)
- Consistent (of course) with fits by Anselmino *et al.*, Vogelsang and Yuan (but different assumptions!)
- Change of sign  $f_{1T}^{\perp a}|_{\text{SIDIS}} = - f_{1T}^{\perp a}|_{\text{DY}}$  can be tested at COMPASS & PAX (byproduct) for Sivers- $q$  ...
- ... also at RHIC, where first glimpse on Sivers- $\bar{q}$  possible
- Most recent & more precise preliminary HERMES data somehow larger Sivers effect. Still compatible with large- $N_c$ .  $\Rightarrow$  even more optimistic estimates for Drell-Yan

*Thank you!*